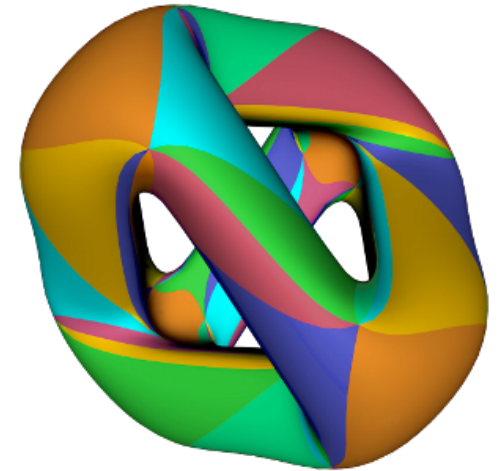
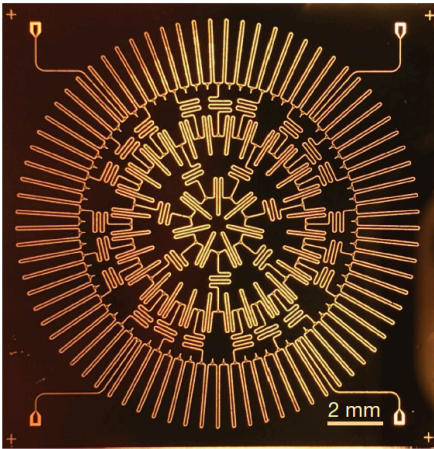


Hyperbolic band theory



Joseph Maciejko
University of Alberta

Discrete LGTs – Emergence & Quantum Simulations

MPQ Garching

May 11, 2022



S. Rayan
(Saskatchewan)

JM & S. Rayan, *Sci. Adv.* 7, eabe9170 (2021)
JM & S. Rayan, *PNAS* 119, e2116869119 (2022)

I. Boettcher, A. V.
Gorshkov, A. J. Kollár, JM,
S. Rayan, R. Thomale,
PRB 105, 125118 (2022)



I. Boettcher
(Alberta)



A. Gorshkov
(JQI/UMD)



A. Kollár
(JQI/UMD)



R. Thomale
(Würzburg)



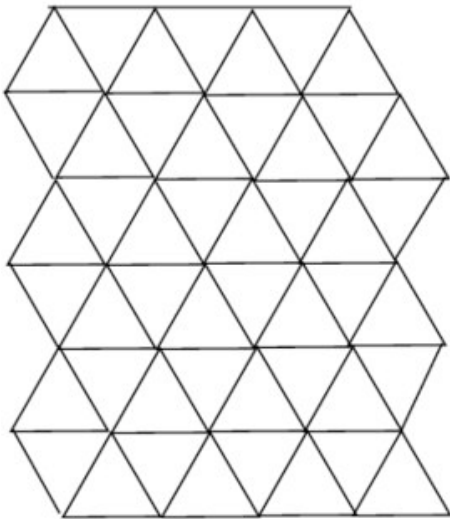
T. Bzdušek
(PSI/Zürich)

T. Bzdušek & JM, in preparation

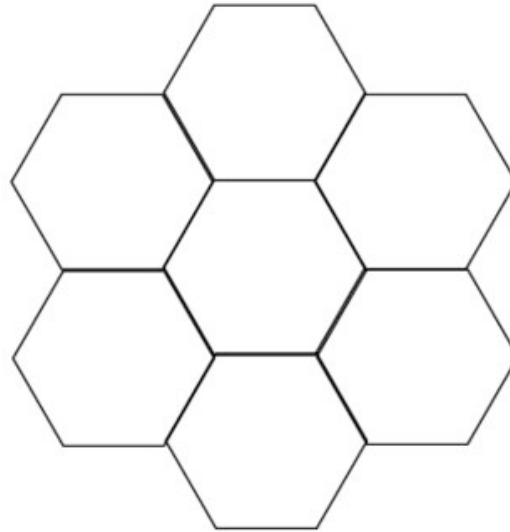
Euclidean lattices



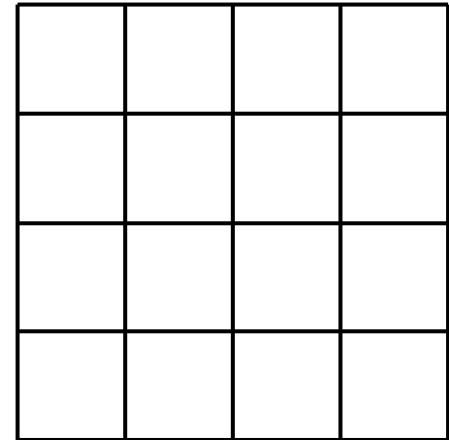
$\{3, 6\}$



$\{6, 3\}$

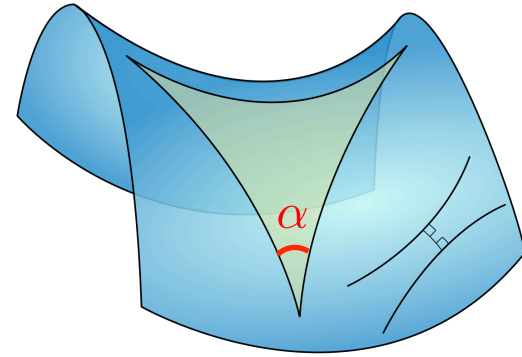


$\{4, 4\}$

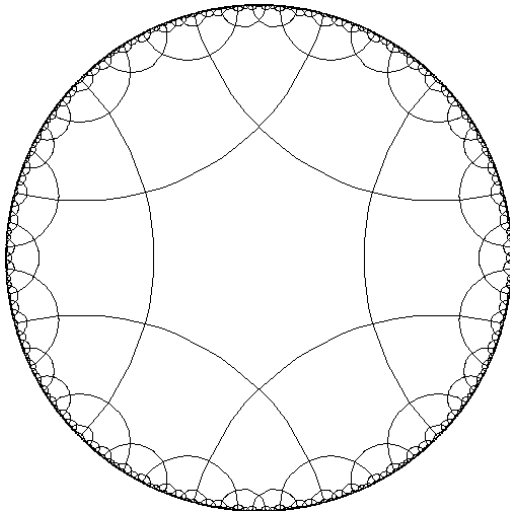


$$(p - 2)(q - 2) = 4$$

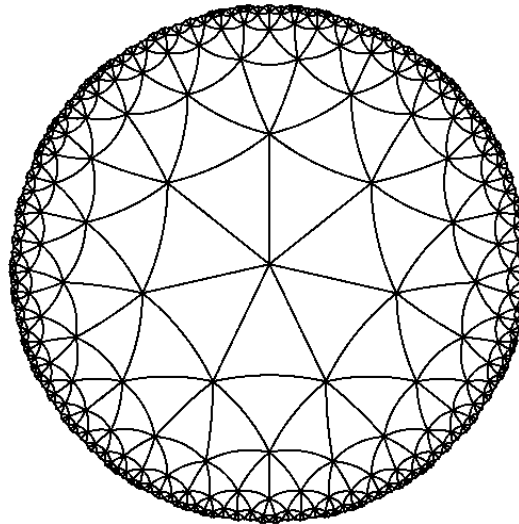
Hyperbolic lattices



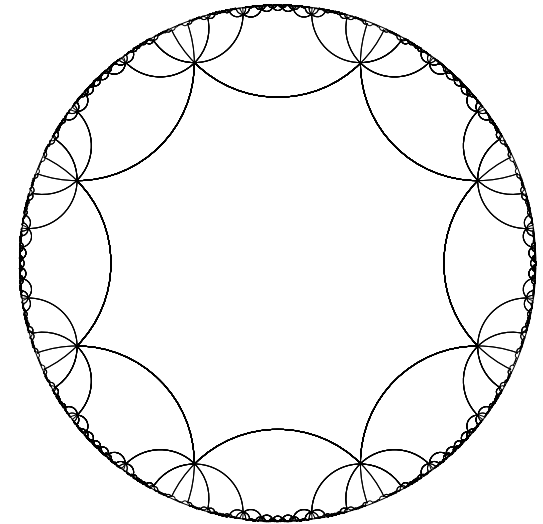
$\{6, 4\}$



$\{3, 7\}$



$\{8, 8\}$



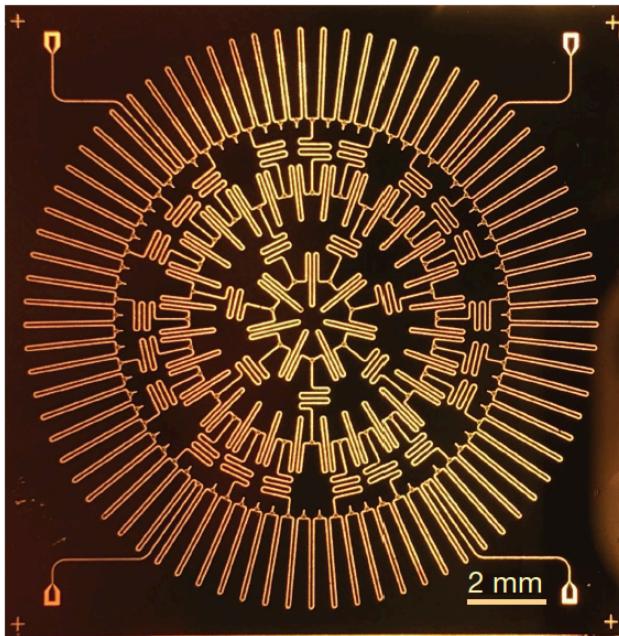
...

$$(p - 2)(q - 2) > 4$$

Hyperbolic lattices in circuit quantum electrodynamics

Alicia J. Kollár^{1,2,3*}, Mattias Fitzpatrick¹ & Andrew A. Houck¹

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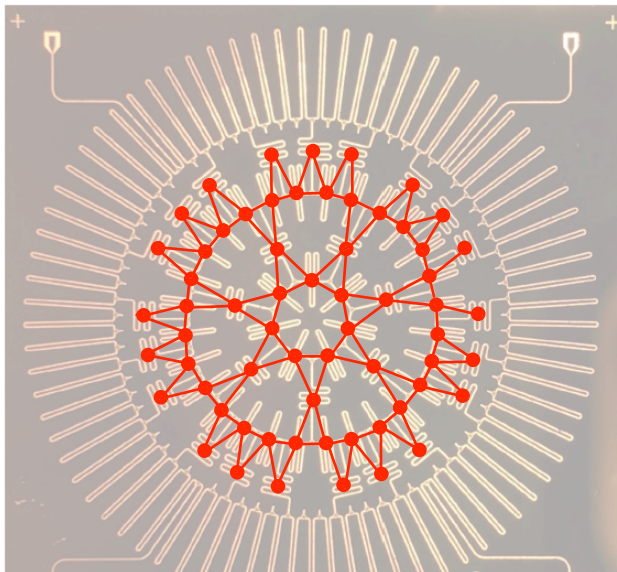


Hyperbolic lattices in circuit quantum electrodynamics

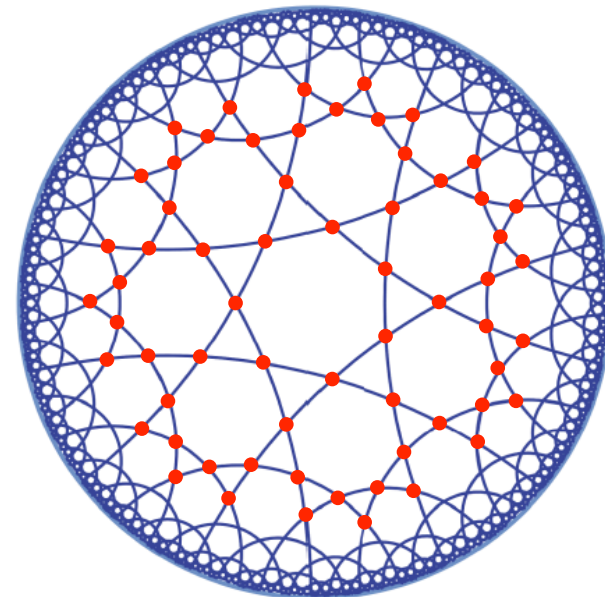
Alicia J. Kollár^{1,2,3*}, Mattias Fitzpatrick¹ & Andrew A. Houck¹

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“hyperbolic kagome lattice”











=

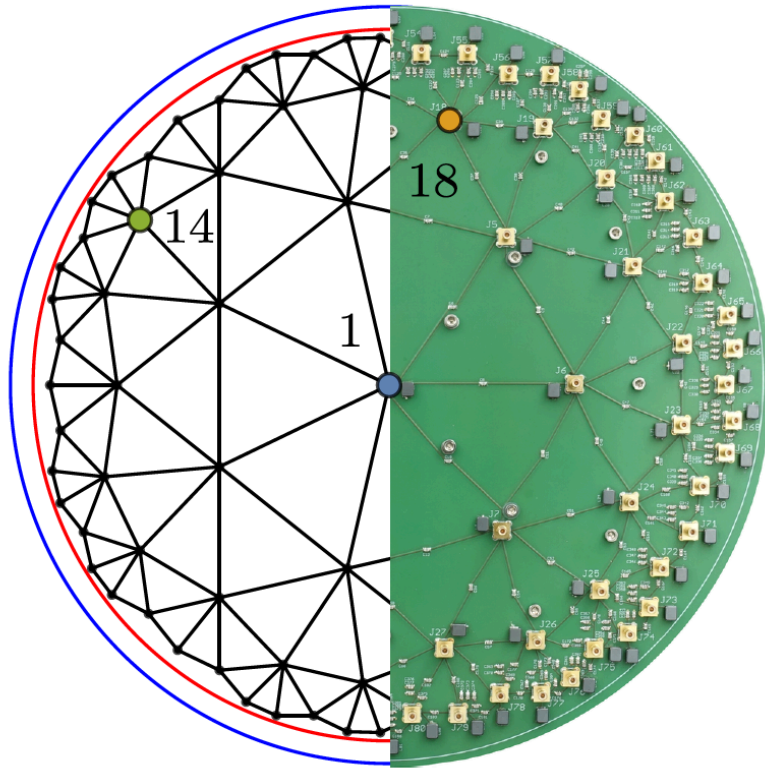


$$H_{\text{TB}} = \omega_0 \sum_i a_i^\dagger a_i - t \sum_{\langle i,j \rangle} (a_i^\dagger a_j + a_j^\dagger a_i)$$

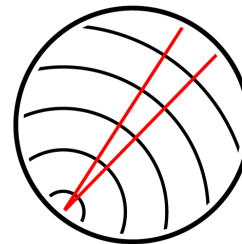
Electric-circuit realization of a hyperbolic drum

Patrick M. Lenggenhager ^{1,2,3,*} Alexander Stegmaier ^{4,*} Lavi K. Upreti ⁴ Tobias Hofmann,⁴ Tobias Helbig ⁴ Achim Vollhardt,² Martin Greiter,⁴ Ching Hua Lee,⁵ Stefan Imhof,⁶ Hauke Brand,⁶ Tobias Kießling,⁶ Igor Boettcher ^{7,8} Titus Neupert ^{2,†} Ronny Thomale ^{4,†} and Tomáš Bzdušek ^{1,2,†}

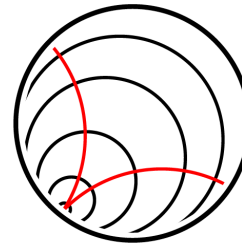
$\{3, 7\}$ “hyperbolic triangular lattice”



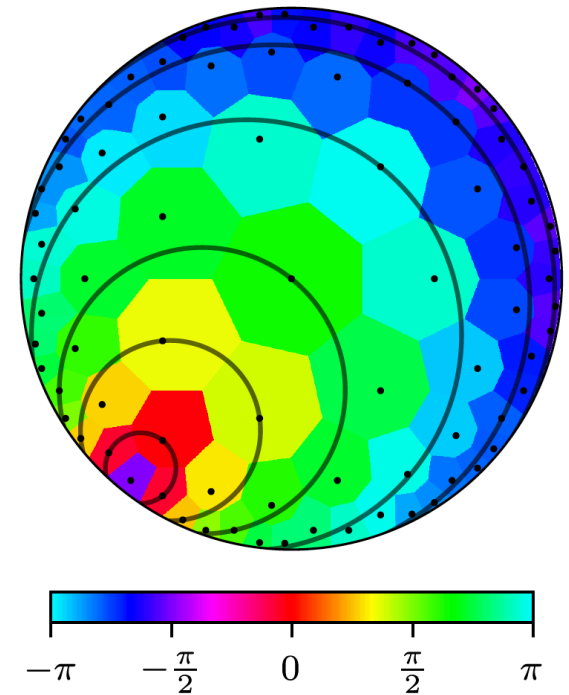
Euclidean drum



Hyperbolic drum



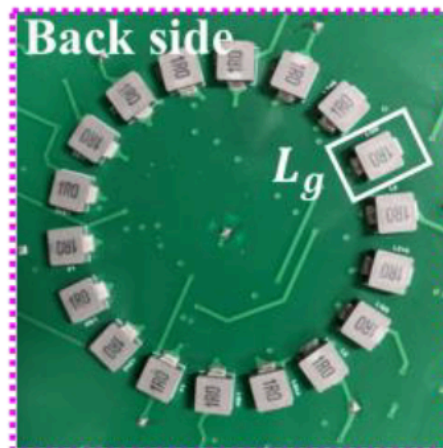
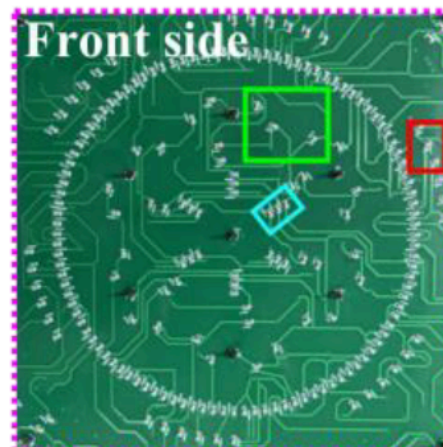
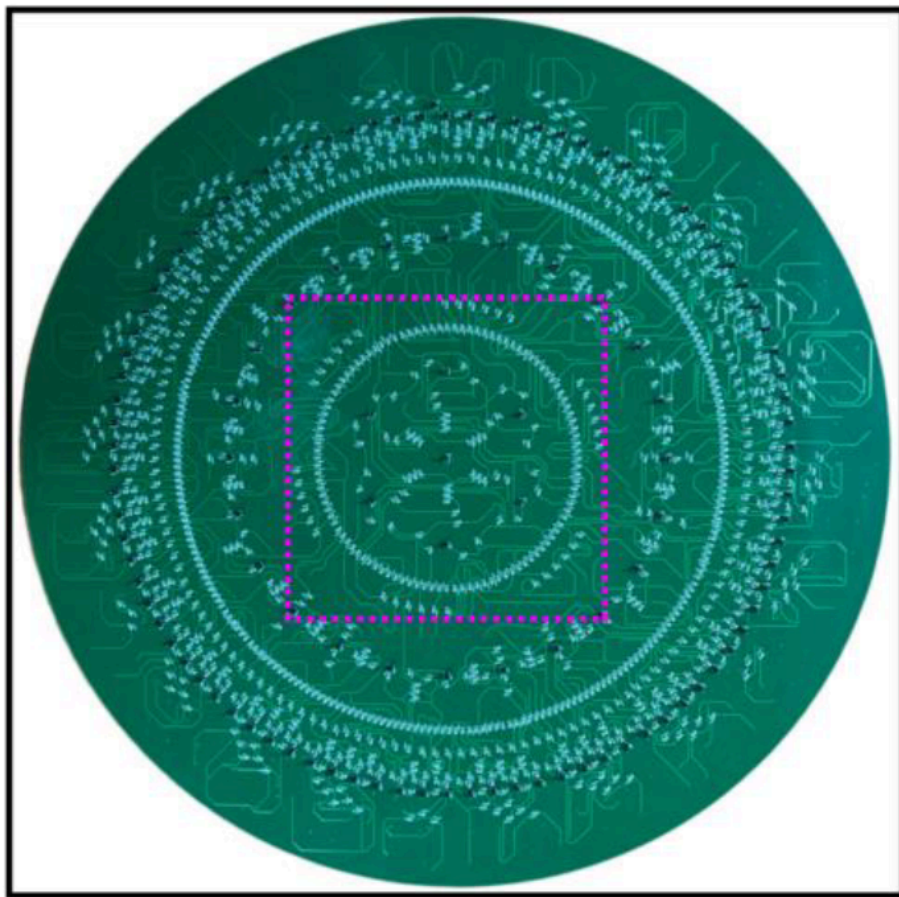
$t = 2.032 \mu\text{s}$



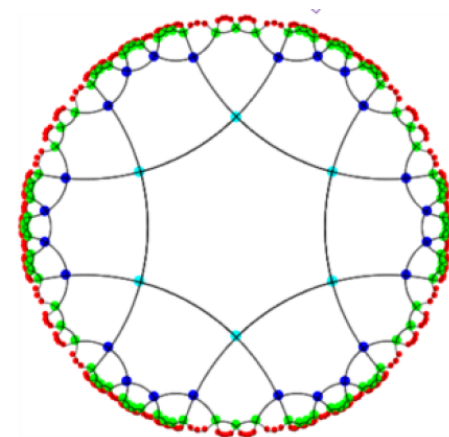
Observation of novel topological states in hyperbolic lattices

Weixuan Zhang^{1*}, Hao Yuan^{1*}, Na Sun¹, Houjun Sun², and Xiangdong Zhang^{1§}

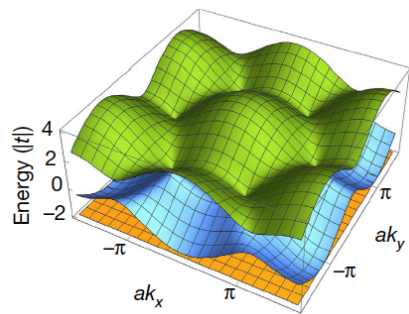
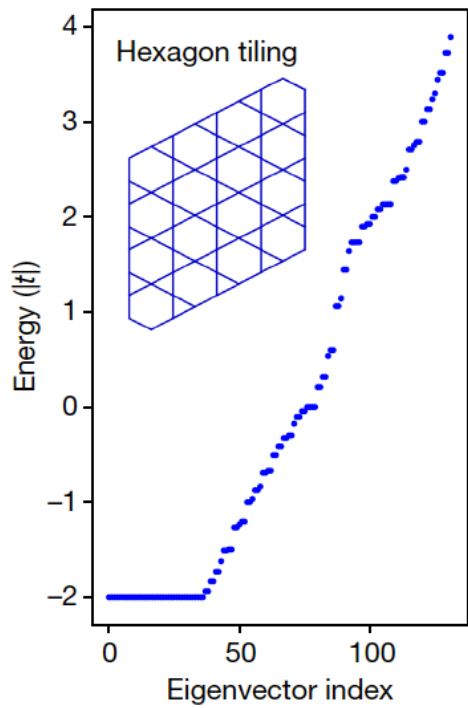
a



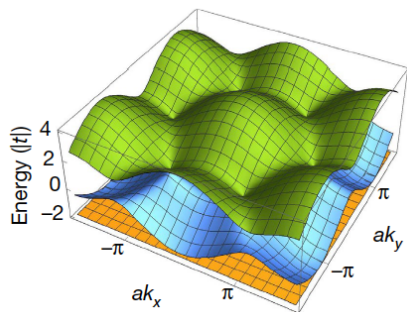
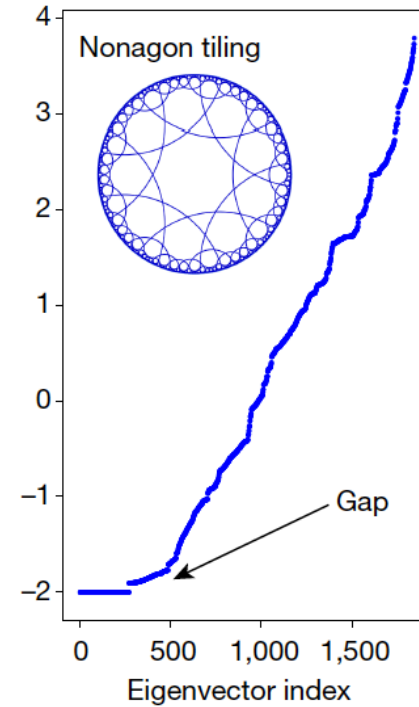
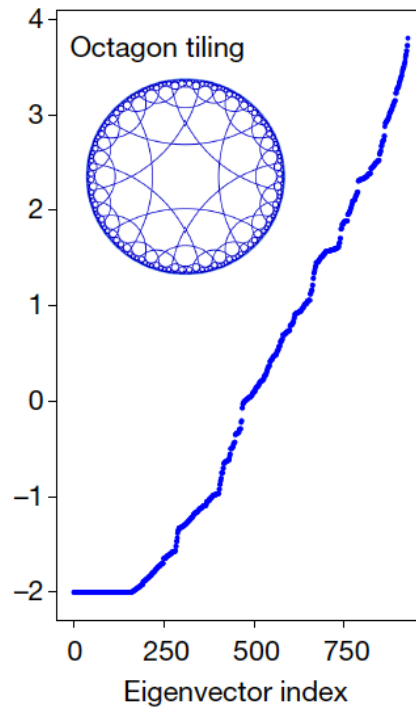
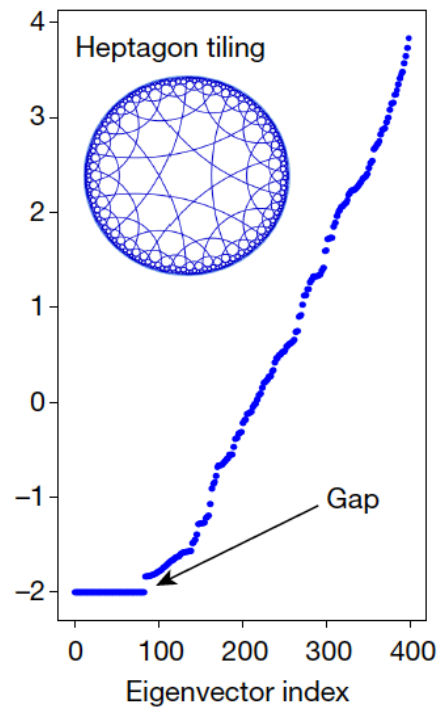
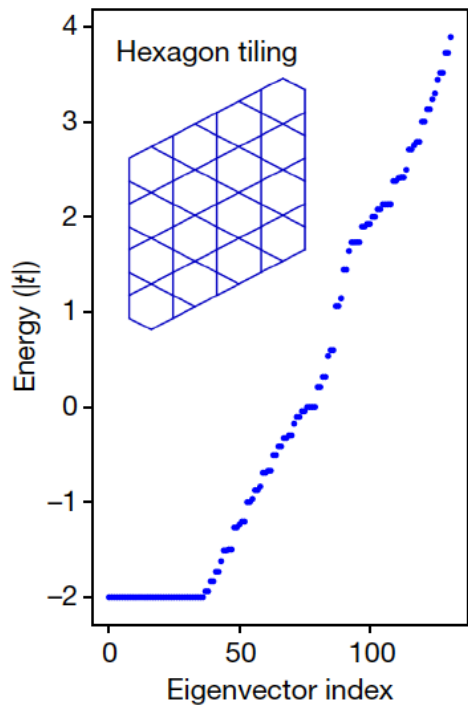
$\{6, 4\}$



“hyperbolic honeycomb lattice”



$E(\mathbf{k})$



$E(k)$



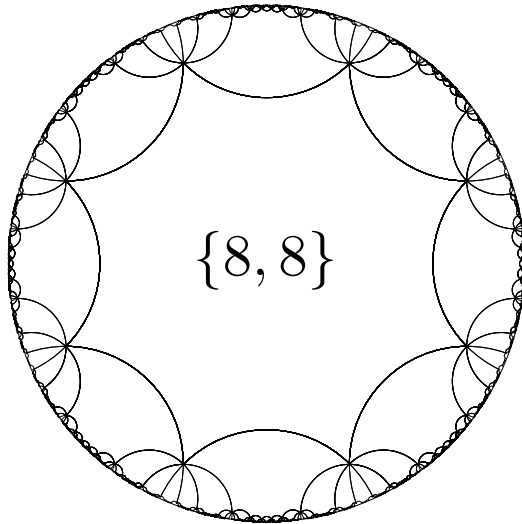
?

Outline

- Hyperbolic geometry & Fuchsian groups
- Hyperbolic Bloch ansatz
- Periodic boundary conditions & automorphic Bloch theorems
- Flat bands & topological bands
- Summary & outlook

Poincaré disk

$$\mathbb{H} = \{|z| < 1\}$$

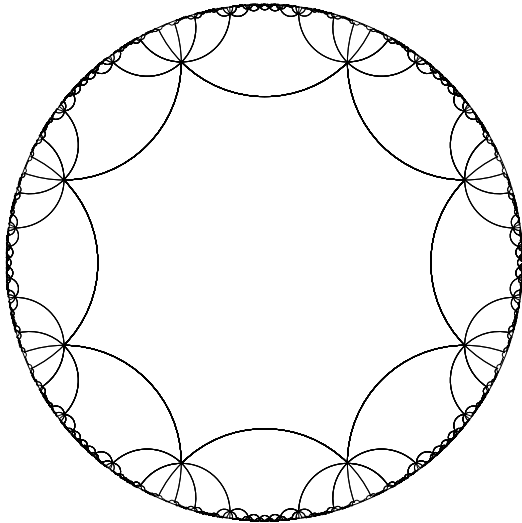


$$ds^2 = \frac{4(dx^2 + dy^2)}{(1 - |z|^2)^2}$$

Poincaré disk



$$\mathbb{H} = \{|z| < 1\}$$



$$ds^2 = \frac{4(dx^2 + dy^2)}{(1 - |z|^2)^2}$$

$\text{PSU}(1,1) \cong \text{PSL}(2,\mathbb{R})$:

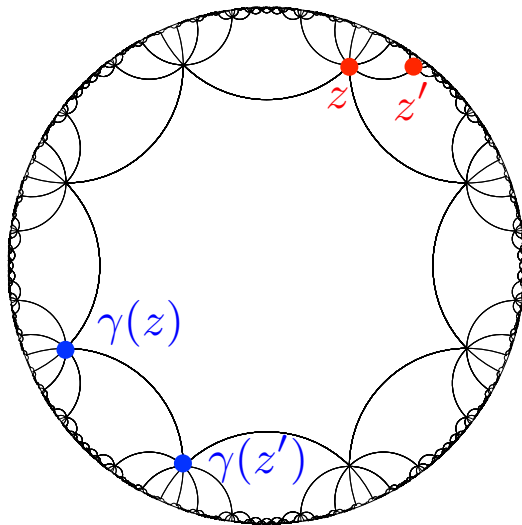
$$\gamma = \begin{pmatrix} \alpha & \beta \\ \beta^* & \alpha^* \end{pmatrix}, \det \gamma = 1$$

$$z \rightarrow \gamma(z) = \frac{\alpha z + \beta}{\beta^* z + \alpha^*}$$

Poincaré disk



$$\mathbb{H} = \{|z| < 1\}$$



$\text{PSU}(1,1) \cong \text{PSL}(2,\mathbb{R})$:

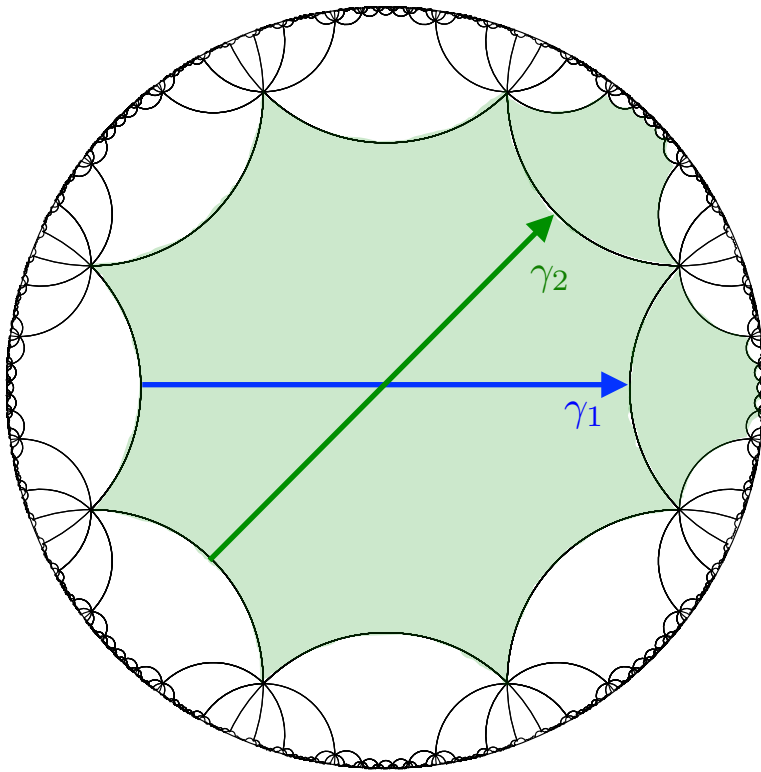
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$$ds^2 = \frac{4(dx^2 + dy^2)}{(1 - |z|^2)^2}$$

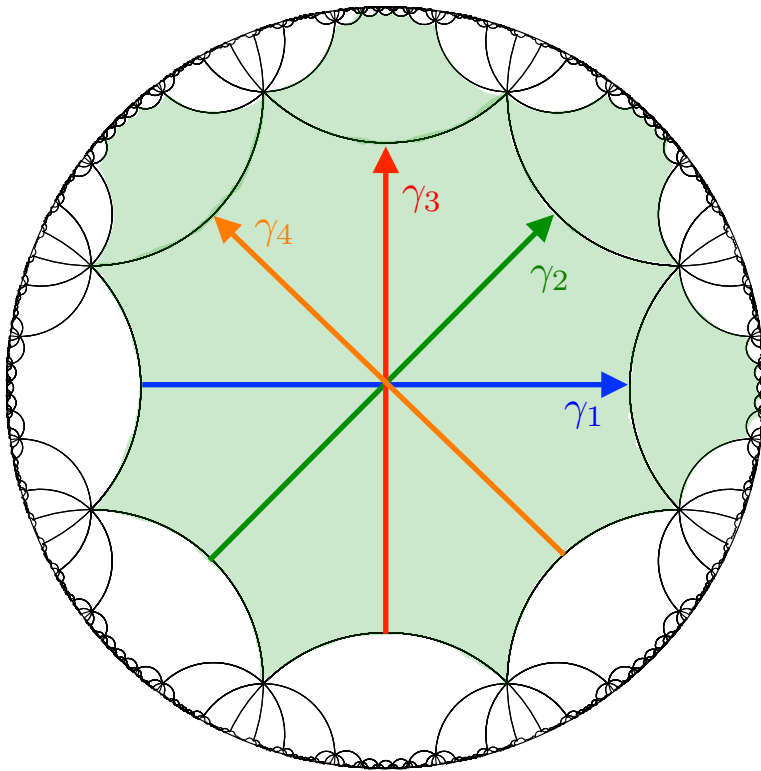
$$d(z, z') = d(\gamma(z), \gamma(z'))$$

Nonabelian translation group



$$\Gamma \subset PSU(1, 1)$$

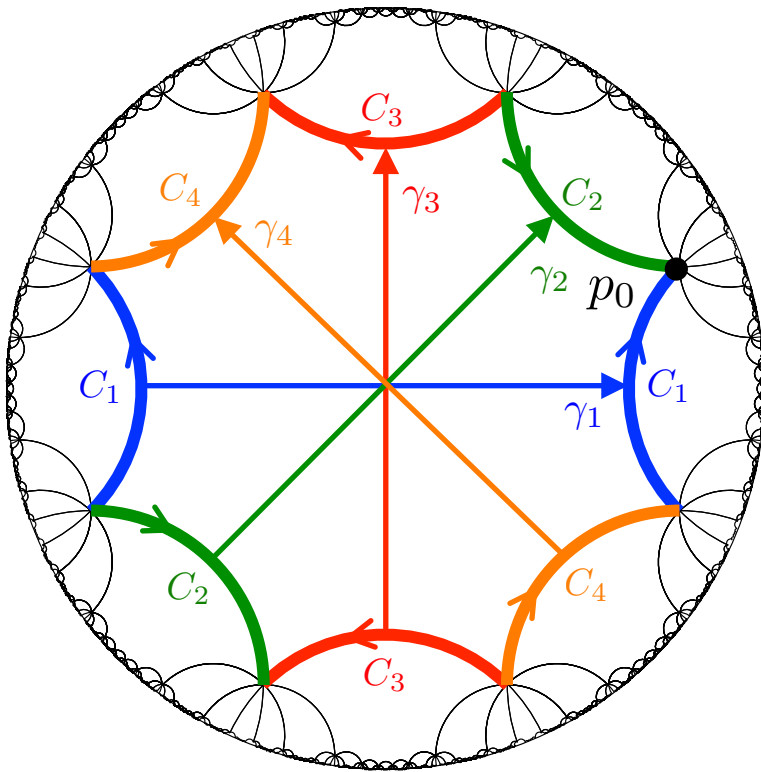
Nonabelian translation group



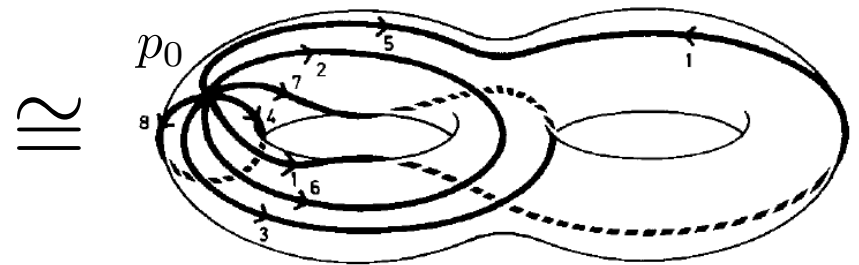
$$\Gamma \subset PSU(1, 1)$$

$$\Gamma = \langle \gamma_1, \gamma_2, \gamma_3, \gamma_4 : \gamma_1 \gamma_2^{-1} \gamma_3 \gamma_4^{-1} \gamma_1^{-1} \gamma_2 \gamma_3^{-1} \gamma_4 = 1 \rangle$$

Compactified unit cell



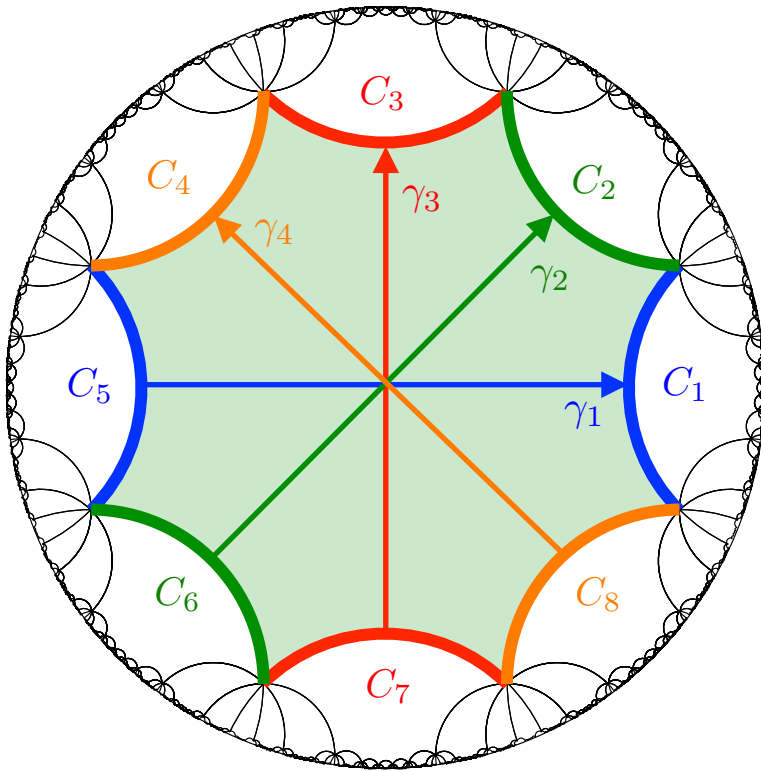
$$\mathbb{H}/\Gamma \cong \Sigma_2$$



$$\Gamma \cong \pi_1(\Sigma_2)$$

$$\gamma_1 \gamma_2^{-1} \gamma_3 \gamma_4^{-1} \gamma_1^{-1} \gamma_2 \gamma_3^{-1} \gamma_4 = 1$$

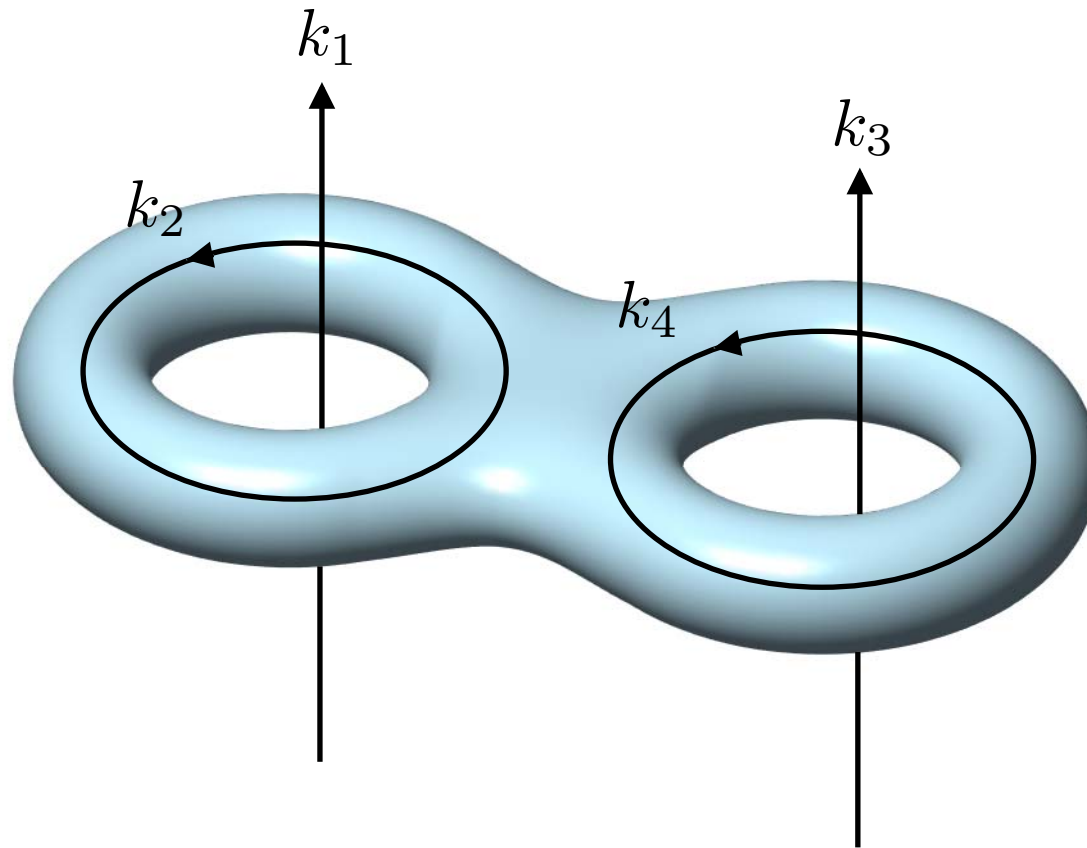
Hyperbolic Bloch ansatz



$$H\psi = E\psi$$

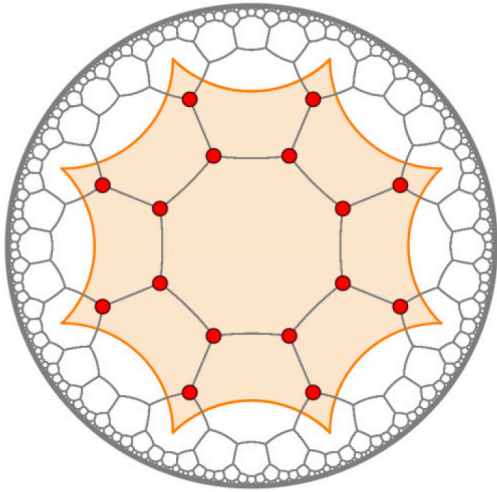
$$\psi(\gamma_j(z)) = e^{ik_j} \psi(z),$$
$$j = 1, 2, 3, 4$$

Hyperbolic crystal momentum

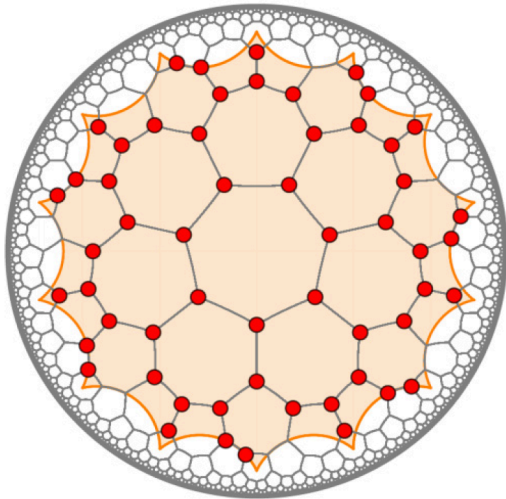


$$\mathbf{k} \equiv (k_1, k_2, \dots, k_{2g-1}, k_{2g}) \in (-\pi, \pi]^{2g} \cong T^{2g} \cong \text{Jac}(\Sigma_g)$$

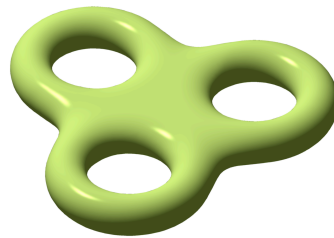
Bravais unit cell



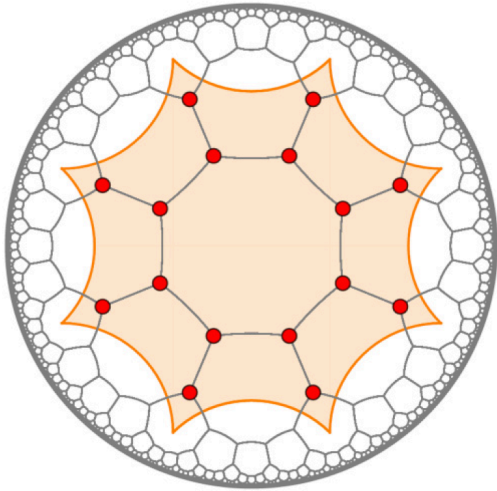
$$\{8,3\} = \{8,8\} + 16\text{-site basis}$$



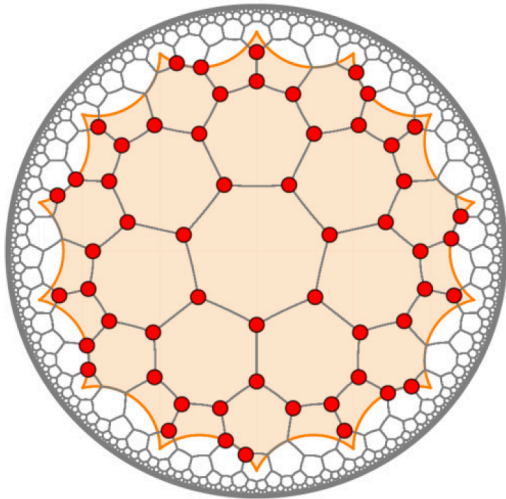
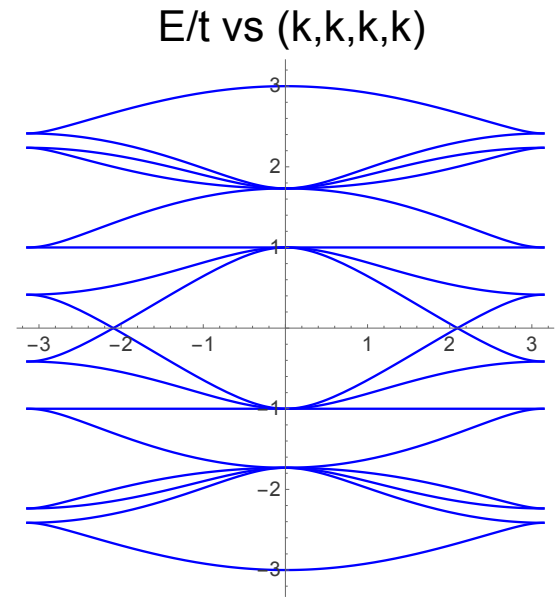
$$\{7,3\} = \{14,7\} + 56\text{-site basis}$$



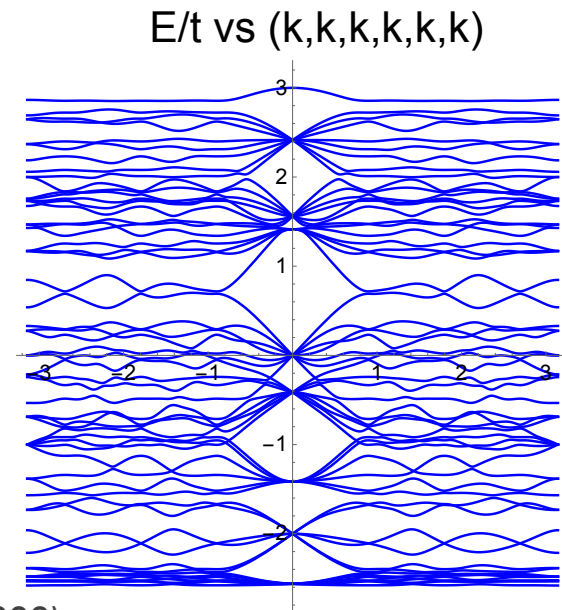
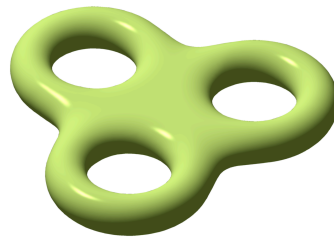
Bravais unit cell



$$\{8,3\} = \{8,8\} + 16\text{-site basis}$$



$$\{7,3\} = \{14,7\} + 56\text{-site basis}$$



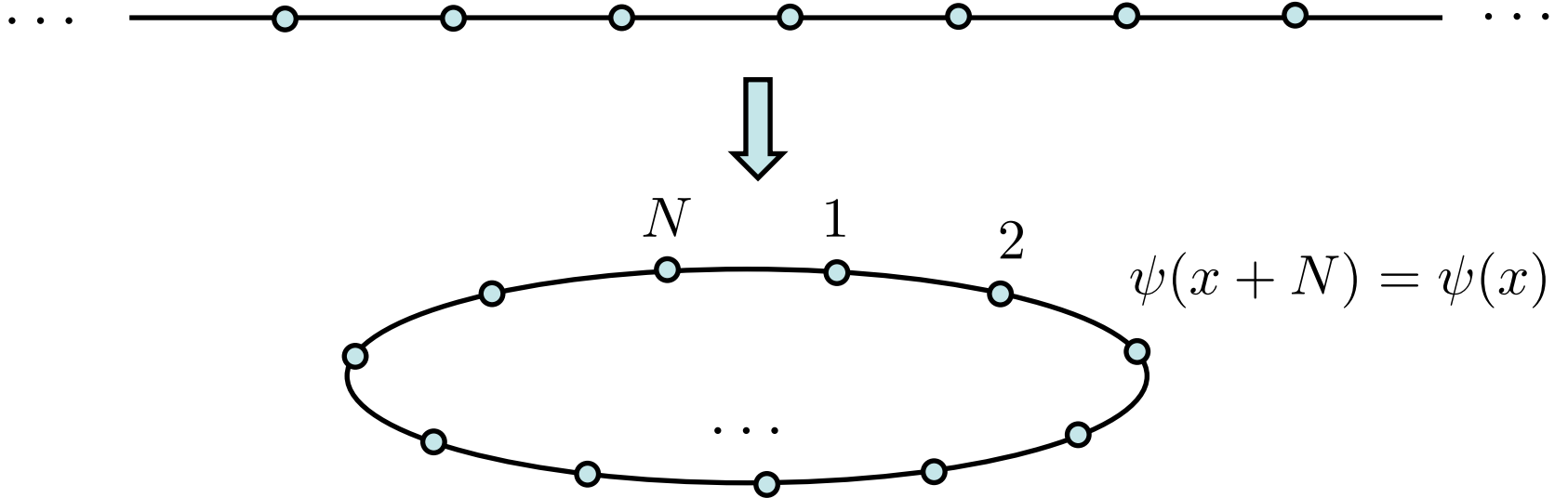
Two issues

- Do hyperbolic Bloch states form a complete set (ansatz vs theorem)?
- What about finite lattices (experiment)?

Two issues

- Do hyperbolic Bloch states form a complete set (ansatz vs theorem)?
- What about finite lattices (experiment)?
- Solution: proper formulation of PBC

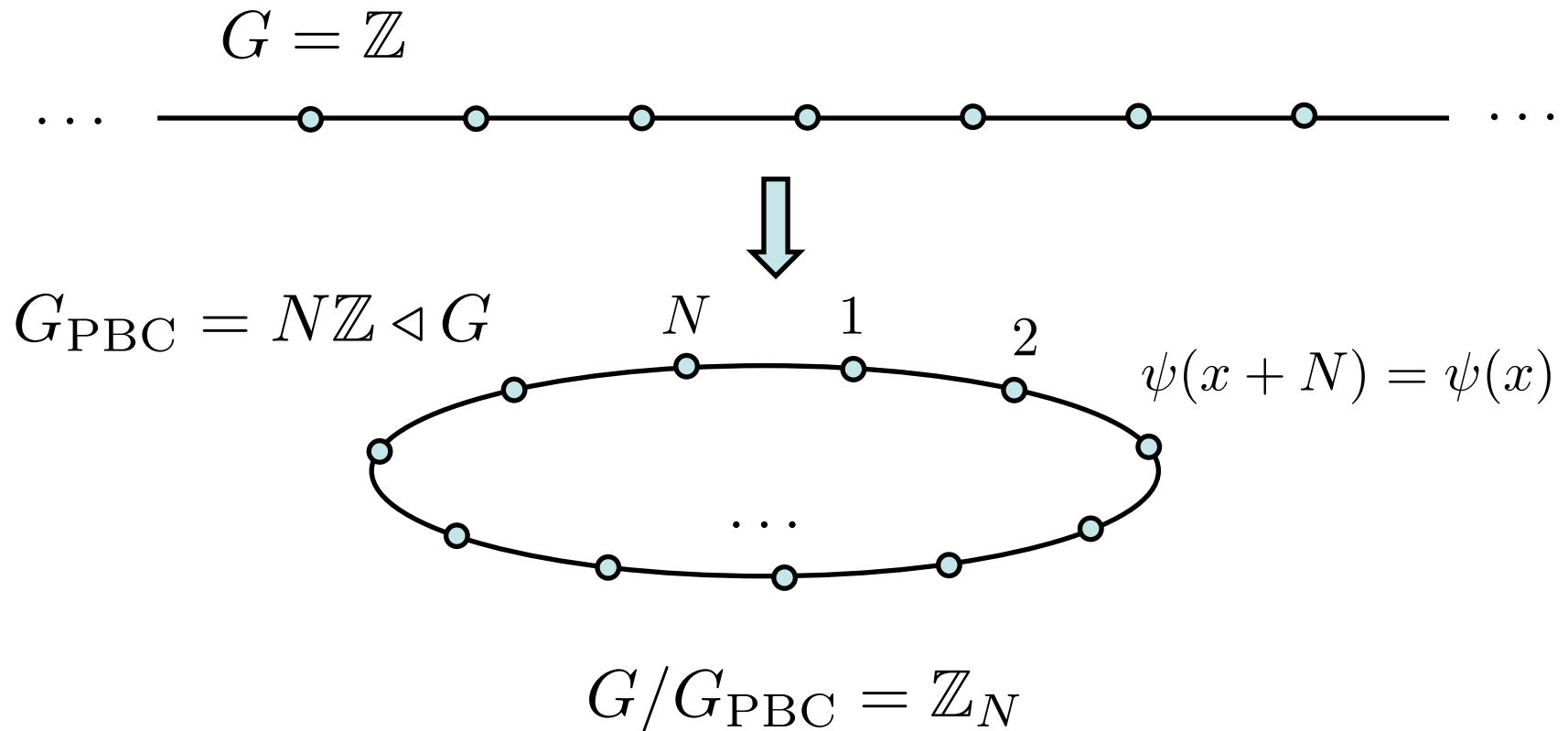
Euclidean PBC



$$k = \frac{2\pi n}{N}, \quad n = 0, 1, \dots, N - 1$$

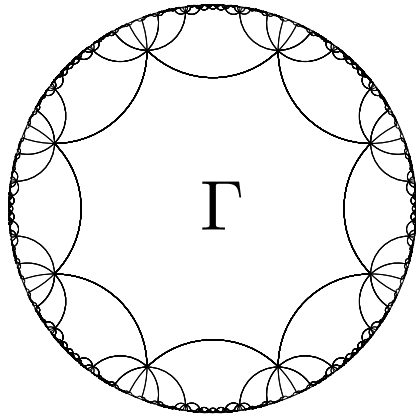
N sites = N Bloch states: complete set

Euclidean PBC: algebraic viewpoint



N allowed k values = N unitary irreps of G/G_{PBC}

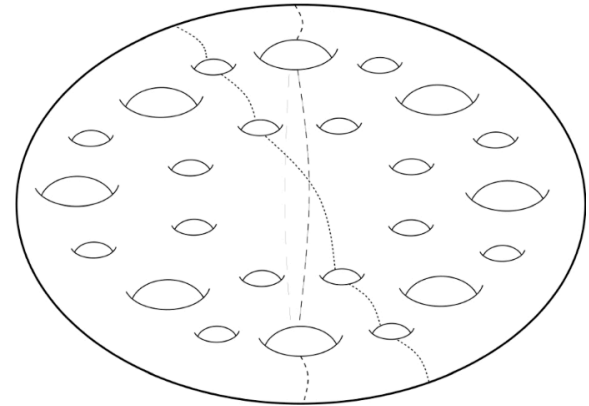
Hyperbolic PBC



infinite lattice



$\Gamma_{\text{PBC}} \triangleleft \Gamma$
normal subgroup of
index N



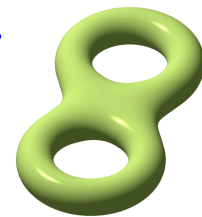
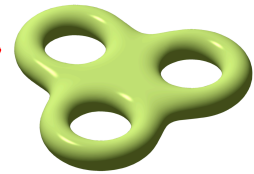
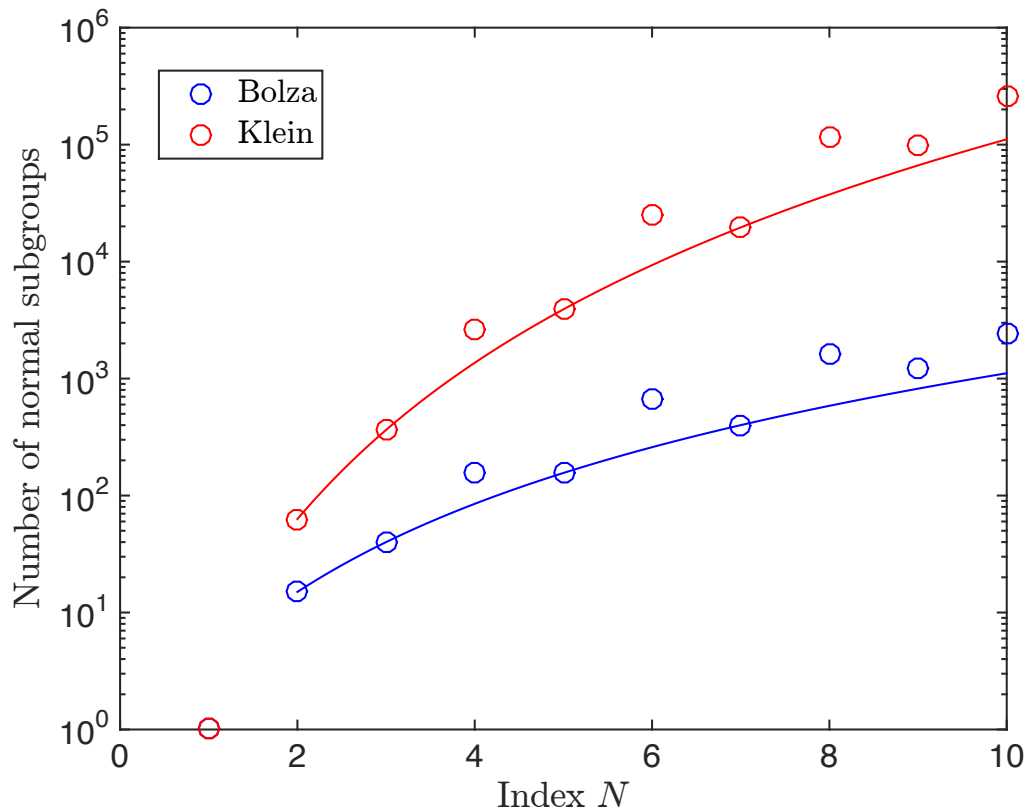
PBC cluster with N unit cells

$$\psi(\gamma_{\text{PBC}}(z)) = \psi(z)$$

Bloch states = unitary irreps of $\Gamma/\Gamma_{\text{PBC}}$

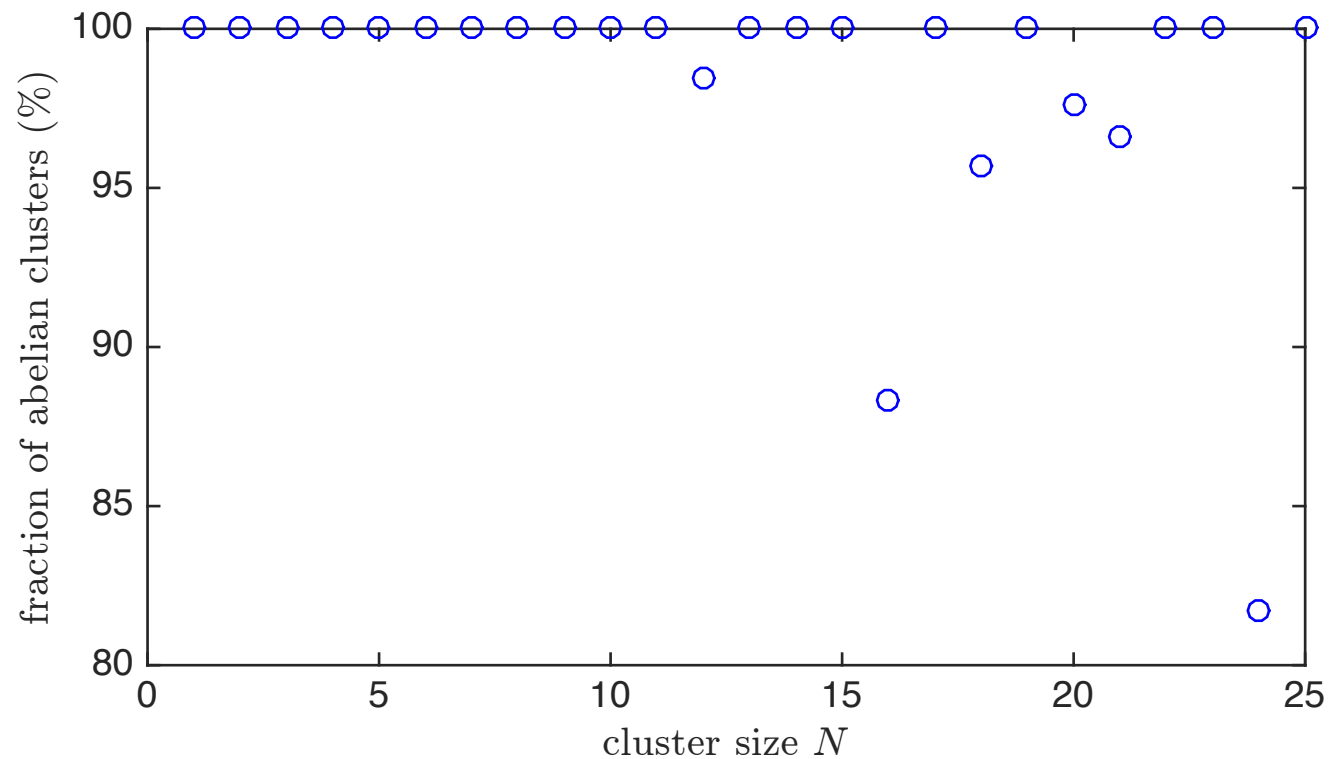
Normal subgroups

- For a given N , many distinct subgroups $\Gamma_{\text{PBC}} \triangleleft \Gamma$
- Enumerate all normal subgroups of index up to N_{max} using computational group theory methods



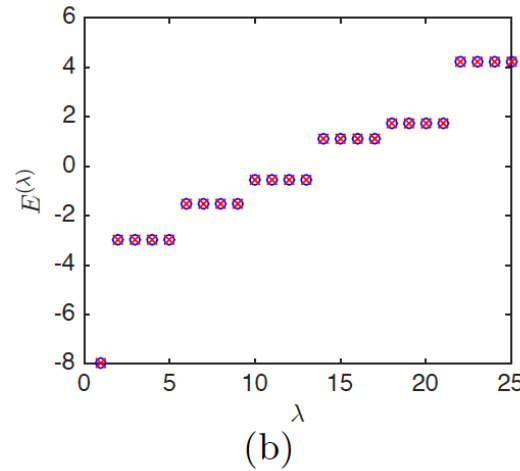
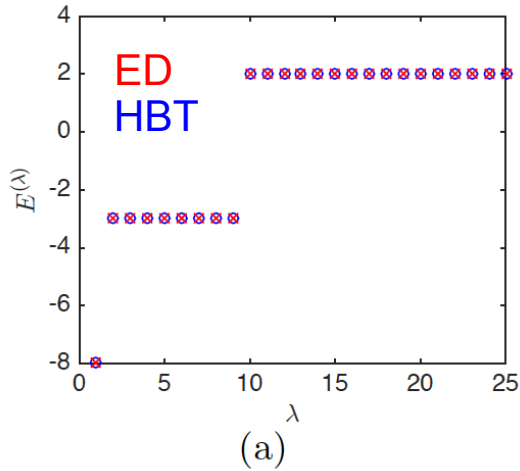
Abelian clusters

- Eigenstates of the clusters fall into irreps of the residual translation group $\Gamma/\Gamma_{\text{PBC}} = \text{finite group of order } N$
- For many clusters (e.g., all prime N), this group is abelian!



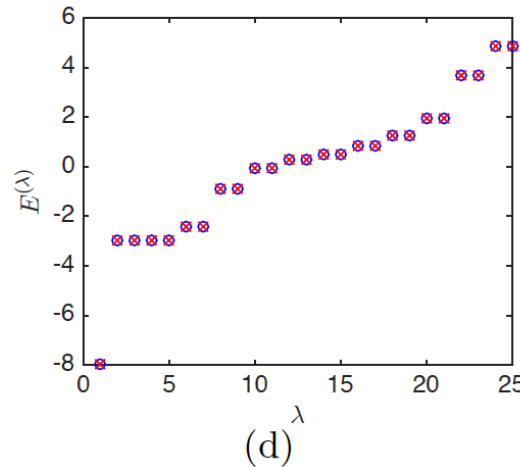
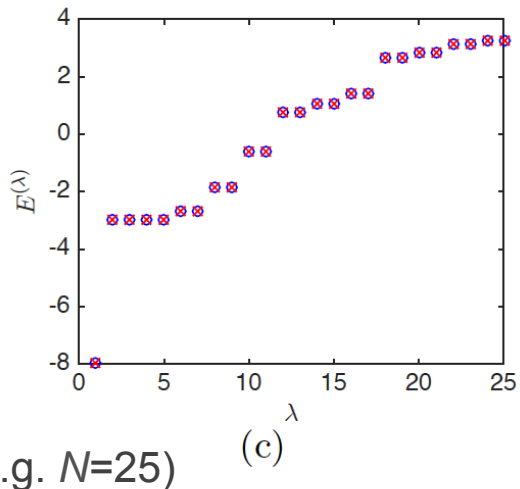
Abelian Bloch theorem

$$\psi^{(\lambda)}(g_k^{-1}(z_i)) = \chi^{(\lambda)}([g_k])\psi^{(\lambda)}(z_i), \quad [g_k] \in \Gamma/\Gamma_{\text{PBC}}$$



$$E^{(\lambda)} = -2 \sum_{j=1}^4 \cos k_j^{(\lambda)},$$

$$\lambda = 1, \dots, N$$



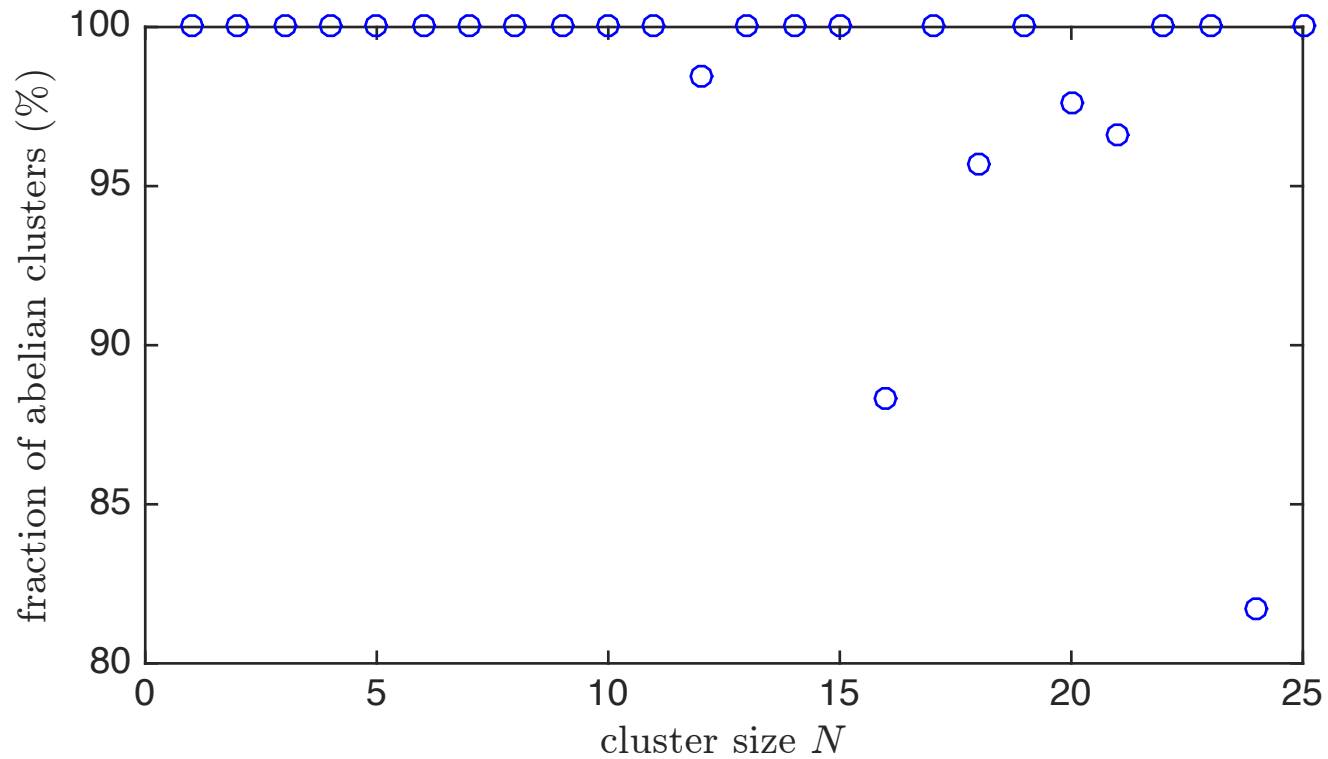
$$k_j^{(\lambda)} \in 2\pi\mathbb{Q}$$



4D BZ is discretized

Nonabelian clusters

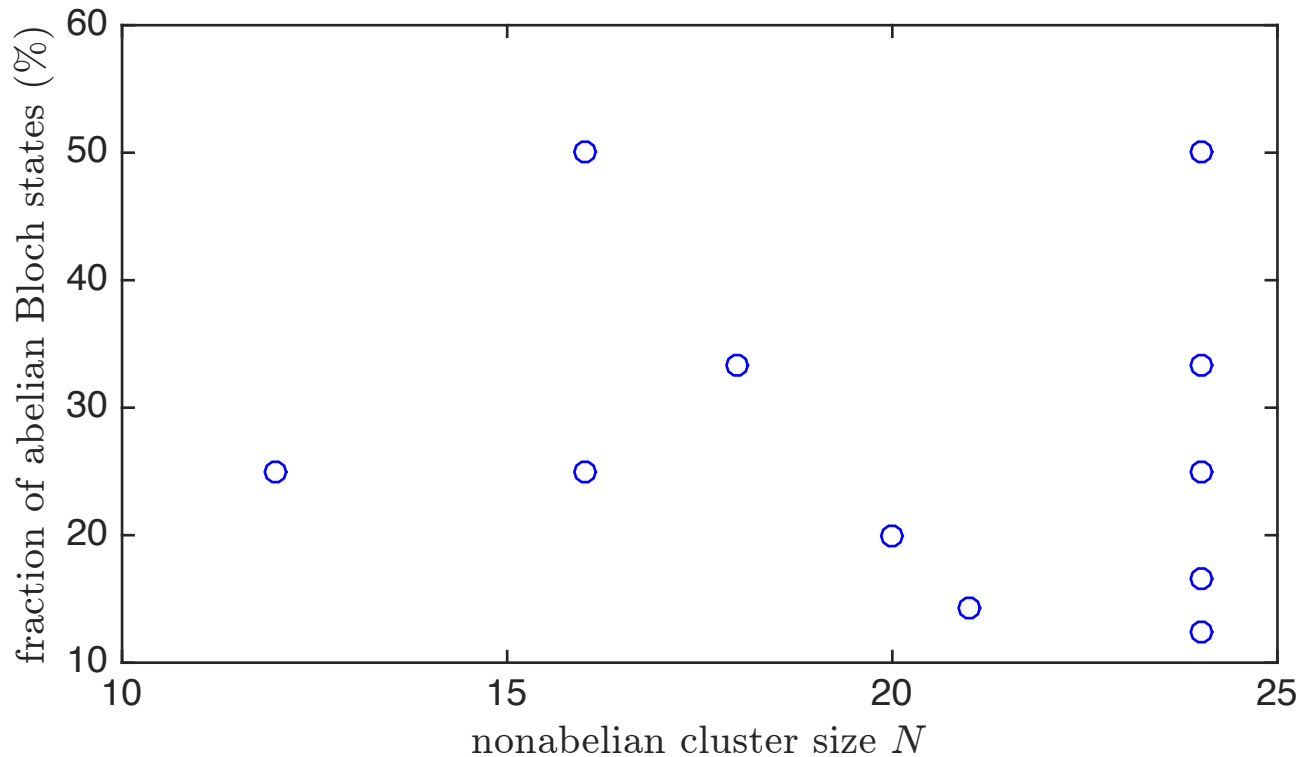
- For $N < 25$, nonabelian $\Gamma/\Gamma_{\text{PBC}}$ found only at $N = 12, 16, 18, 20, 21, 24$



Nonabelian Bloch theorem

- Nonabelian $\Gamma/\Gamma_{\text{PBC}}$ possesses higher-dimensional unitary irreps:

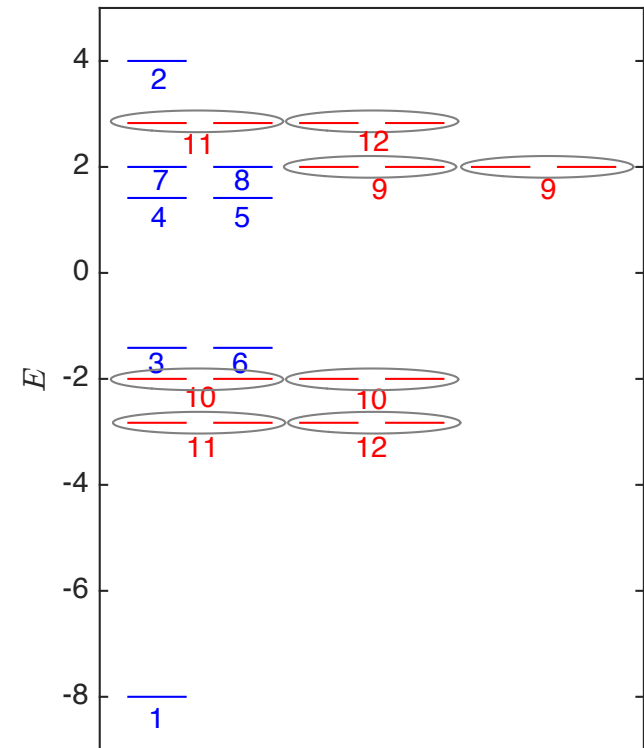
$$\psi_{\nu}^{(\lambda)}(g_k^{-1}(z_i)) = \sum_{\mu=1}^{r_{\lambda}} \psi_{\mu}^{(\lambda)}(z_i) D_{\mu\nu}^{(\lambda)}([g_k]), \quad [g_k] \in \Gamma/\Gamma_{\text{PBC}}$$



Nonabelian Bloch theorem

- Example ($N=24$): 8 abelian irreps, 4 nonabelian (2D) irreps

C	1	2	3	4	5	6	7	8	9	10	11	12
n_C	1	1	1	1	2	2	2	2	3	3	3	3
$D^{(1)}$	1	1	1	1	1	1	1	1	1	1	1	1
$D^{(2)}$	1	1	1	1	1	1	1	1	-1	-1	-1	-1
$D^{(3)}$	1	-1	a	$-a$	1	-1	a	$-a$	c	$-c$	$-1/c$	$1/c$
$D^{(4)}$	1	-1	a	$-a$	1	-1	a	$-a$	$-c$	c	$1/c$	$-1/c$
$D^{(5)}$	1	-1	$-a$	a	1	-1	$-a$	a	$-1/c$	$1/c$	c	$-c$
$D^{(6)}$	1	-1	$-a$	a	1	-1	$-a$	a	$1/c$	$-1/c$	$-c$	c
$D^{(7)}$	1	1	-1	-1	1	1	-1	-1	a	a	$-a$	$-a$
$D^{(8)}$	1	1	-1	-1	1	1	-1	-1	$-a$	$-a$	a	a
$D^{(9)}$	2	2	-2	-2	-1	-1	1	1	0	0	0	0
$D^{(10)}$	2	2	2	2	-1	-1	-1	-1	0	0	0	0
$D^{(11)}$	2	-2	b	$-b$	-1	1	$-a$	a	0	0	0	0
$D^{(12)}$	2	-2	$-b$	b	-1	1	a	$-a$	0	0	0	0



Nonabelian Brillouin zones

$$\text{Irrep}(\Gamma, U(r))/U(r)$$

(topological)



Narasimhan-Seshadri
(1965)

Riemann-Hilbert

flat $U(r)$ connections on Σ_g
gauge group

(smooth)

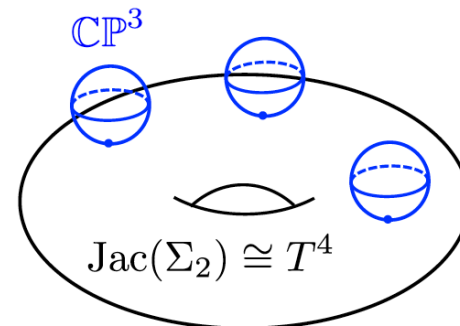
Atiyah-Bott, Witten...

$$\mathcal{M}(\Sigma_g, U(r))$$

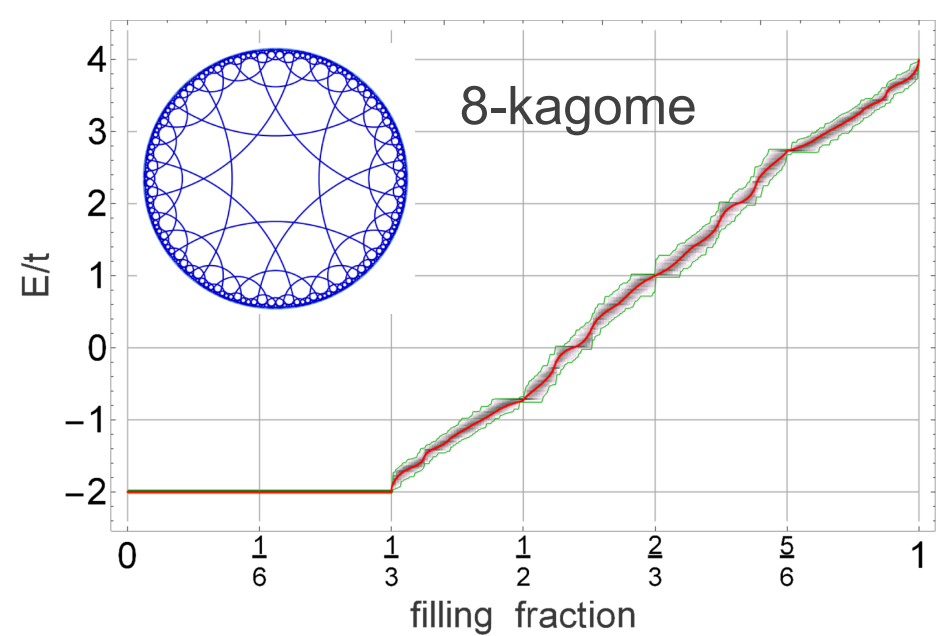
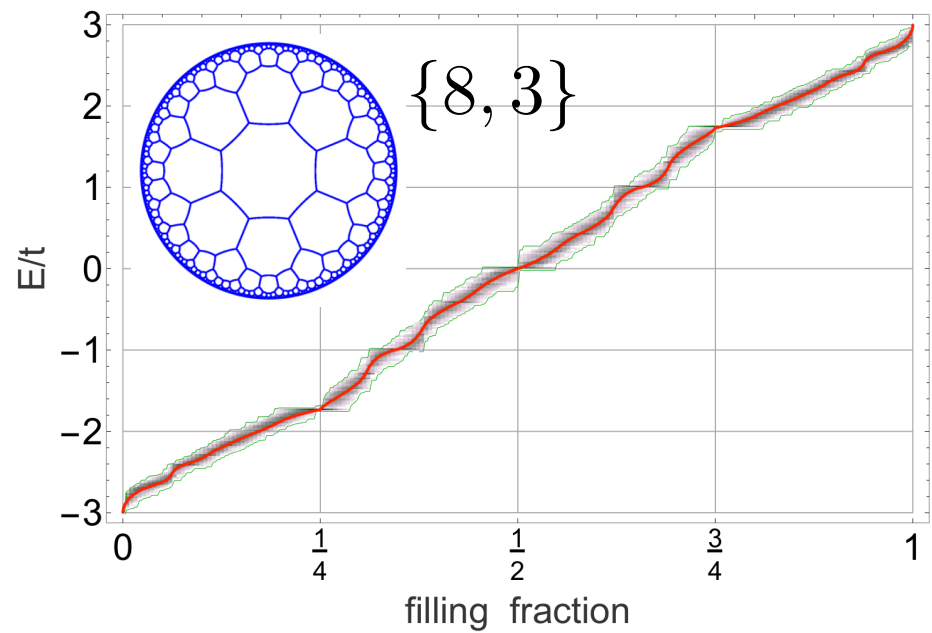
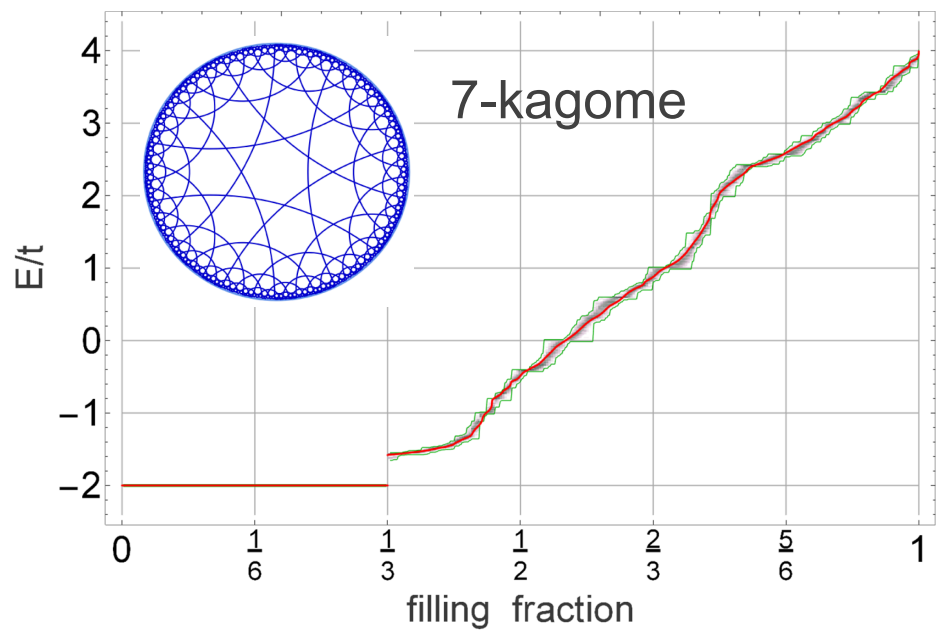
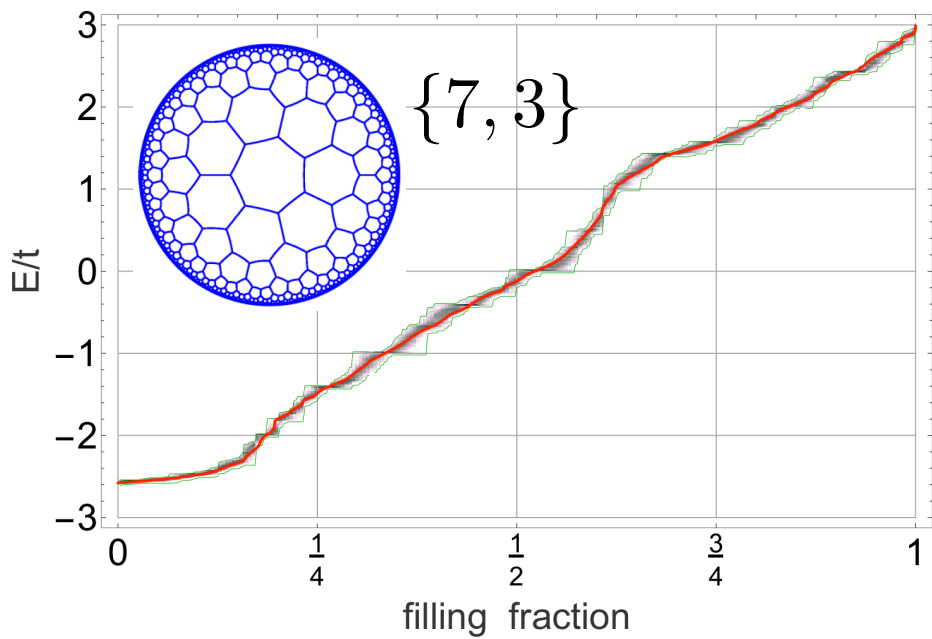
moduli space of stable holomorphic
vector bundles of rank r

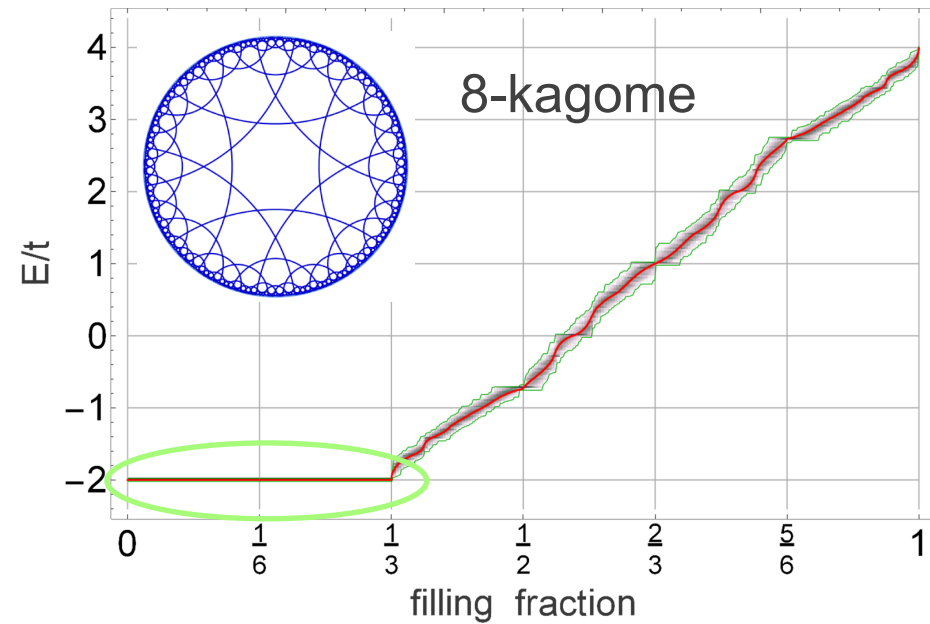
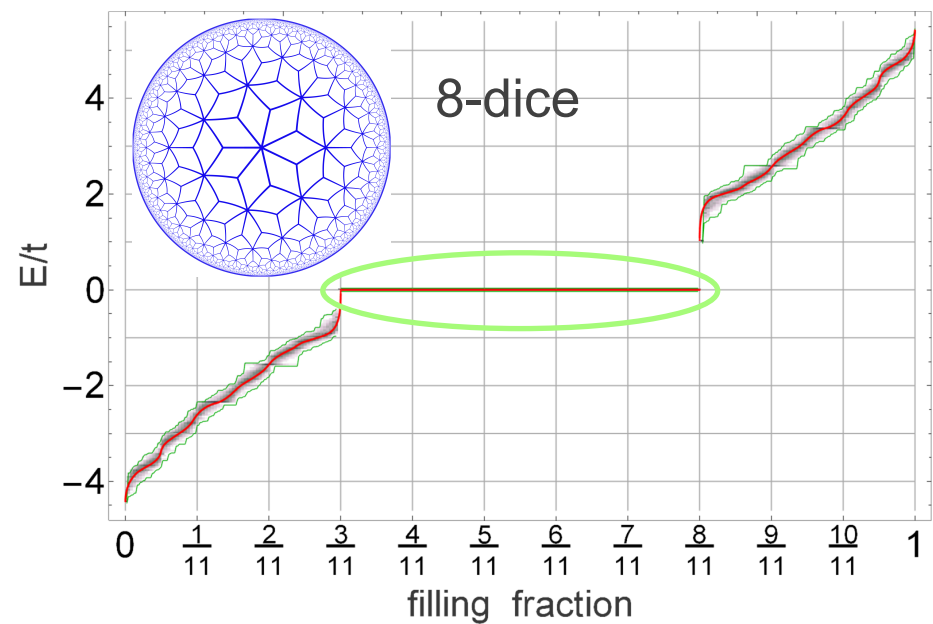
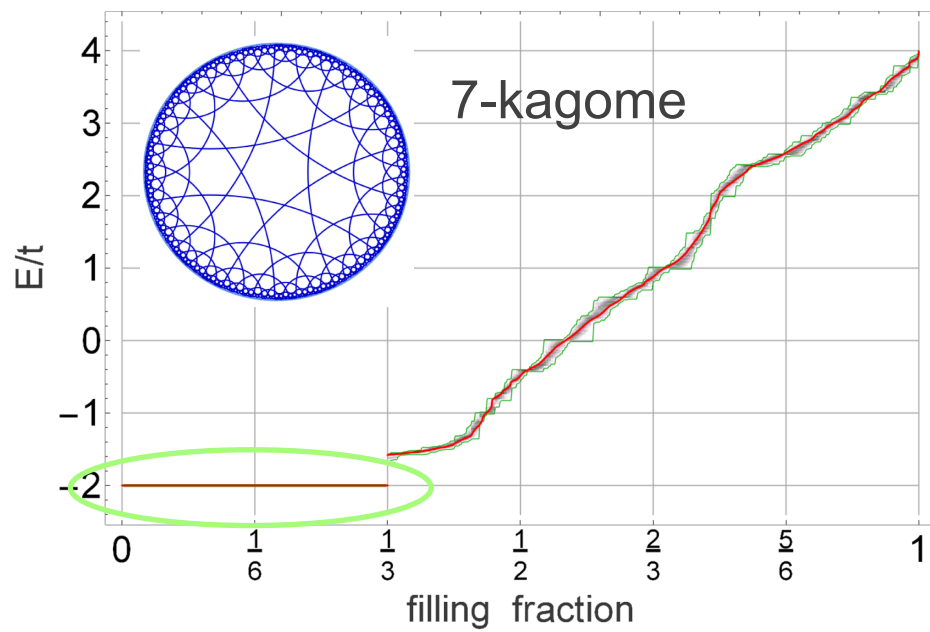
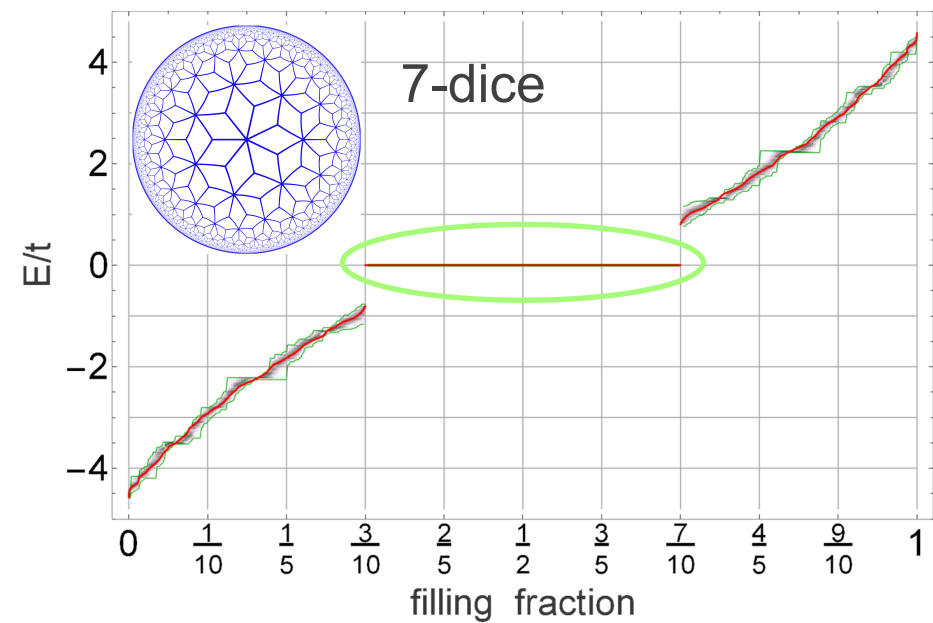
(holomorphic)

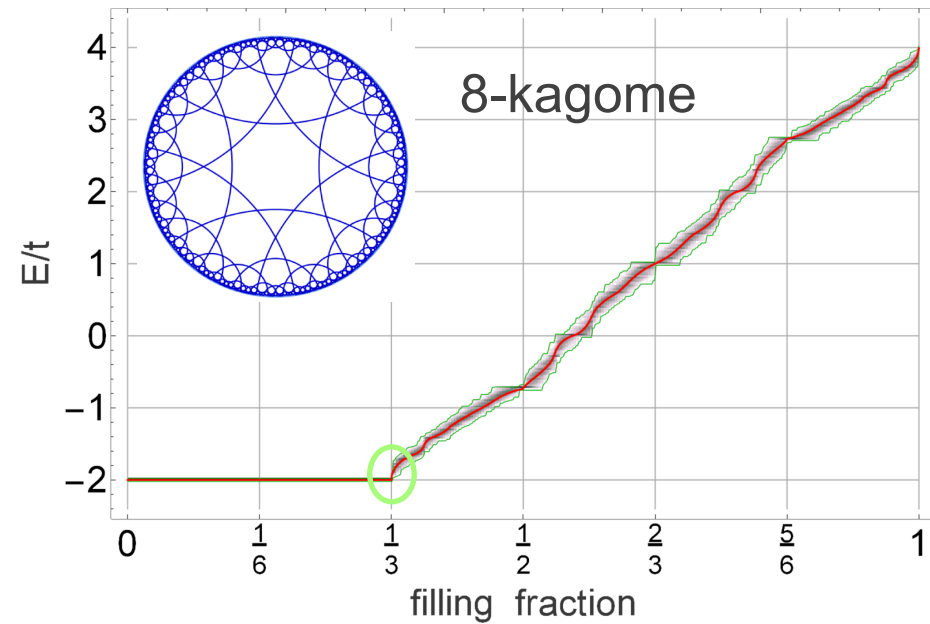
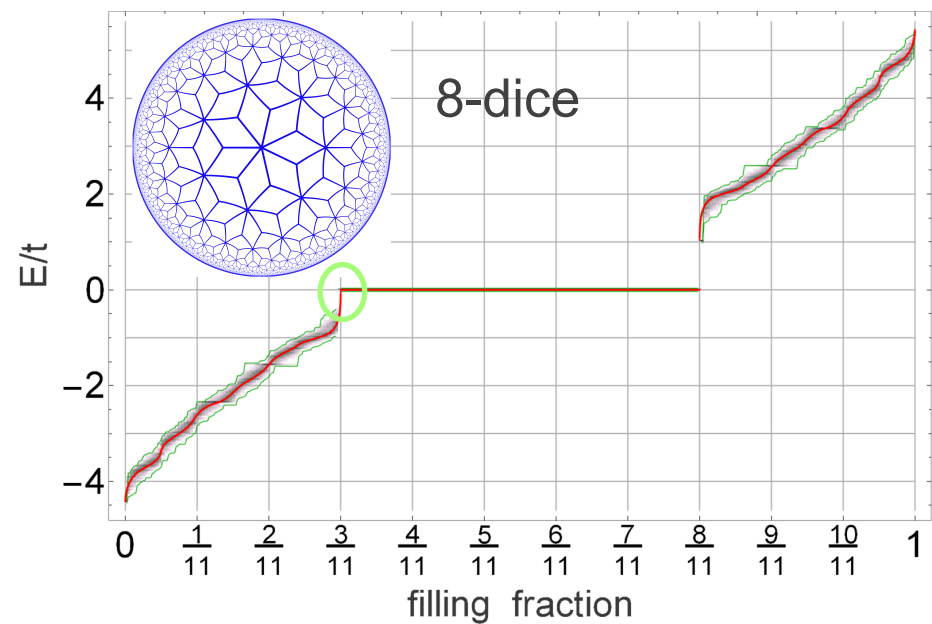
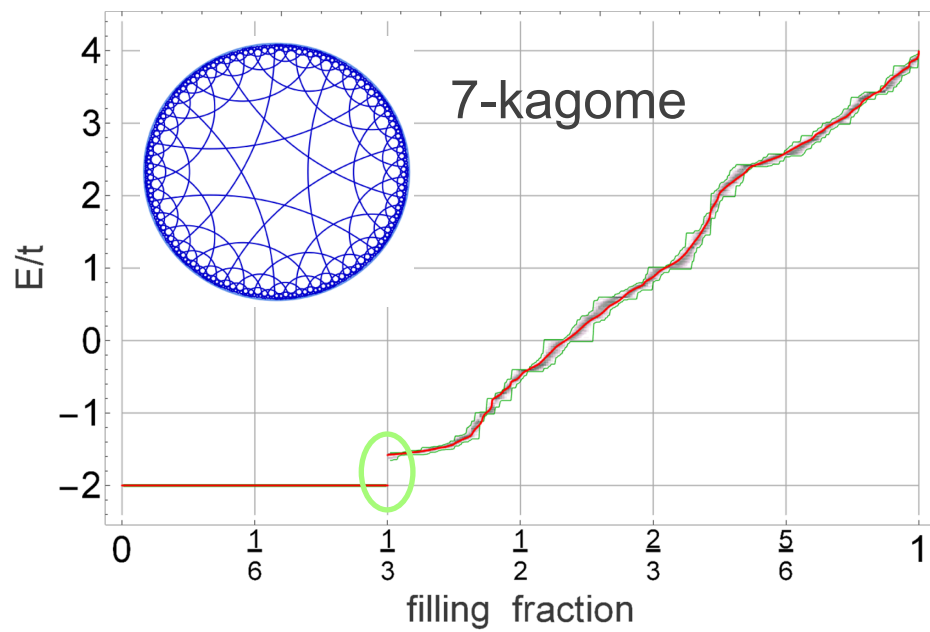
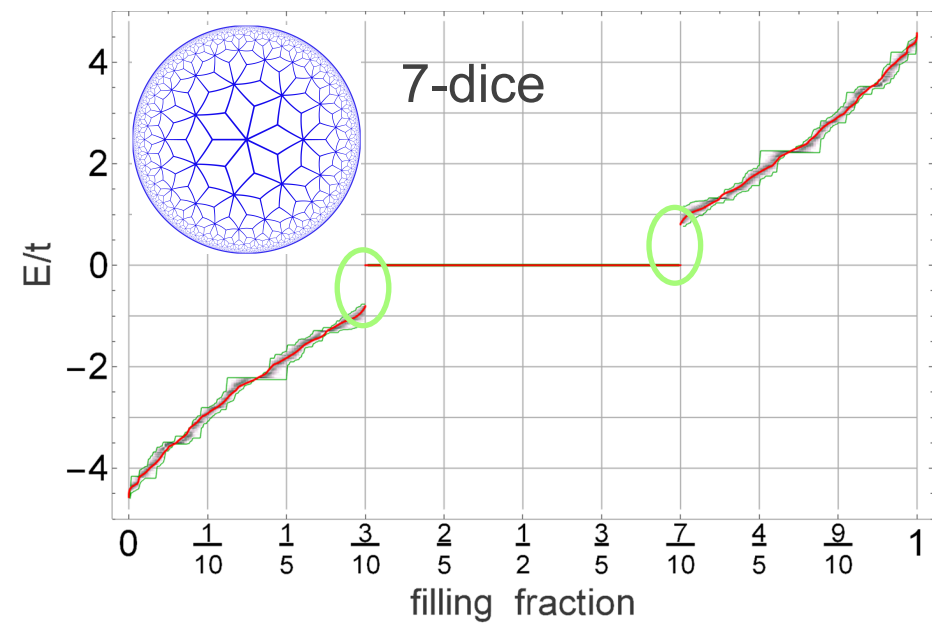
{8,8} lattice, 2D irreps: $\mathcal{M} \approx T^4 \times \mathbb{C}P^3$



Are higher-dimensional irreps important?







Hyperbolic band topology?

Chern insulator in a hyperbolic lattice

Zheng-Rong Liu,¹ Chun-Bo Hua,¹ Tan Peng,¹ and Bin Zhou^{1,*}






arXiv:2203.02101 (th)

Observation of novel topological states in hyperbolic lattices

Weixuan Zhang^{1*}, Hao Yuan^{1*}, Na Sun¹, Houjun Sun², and Xiangdong Zhang^{1§}

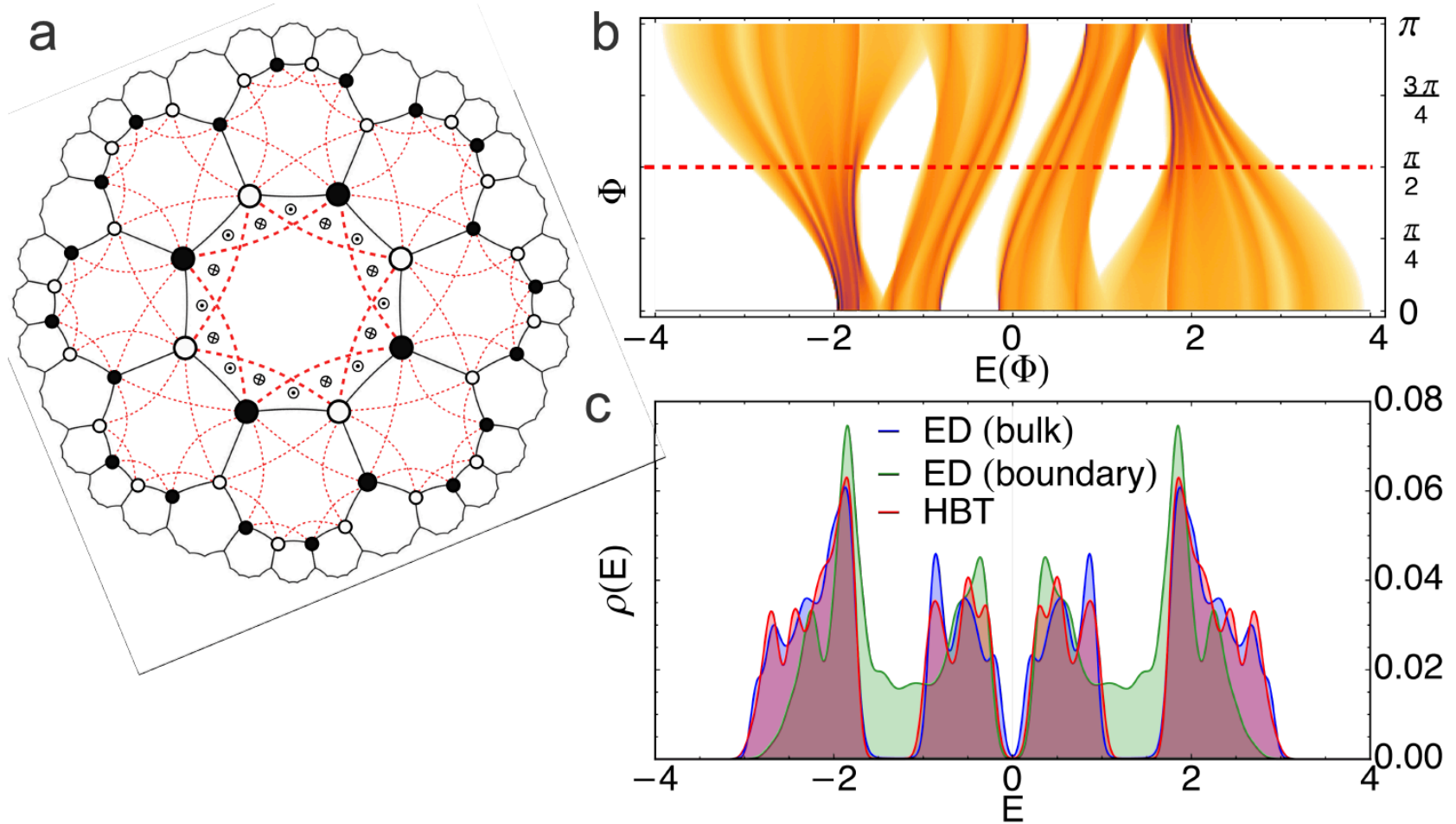
arXiv:2203.03214 (th+exp)

Hyperbolic topological band insulators

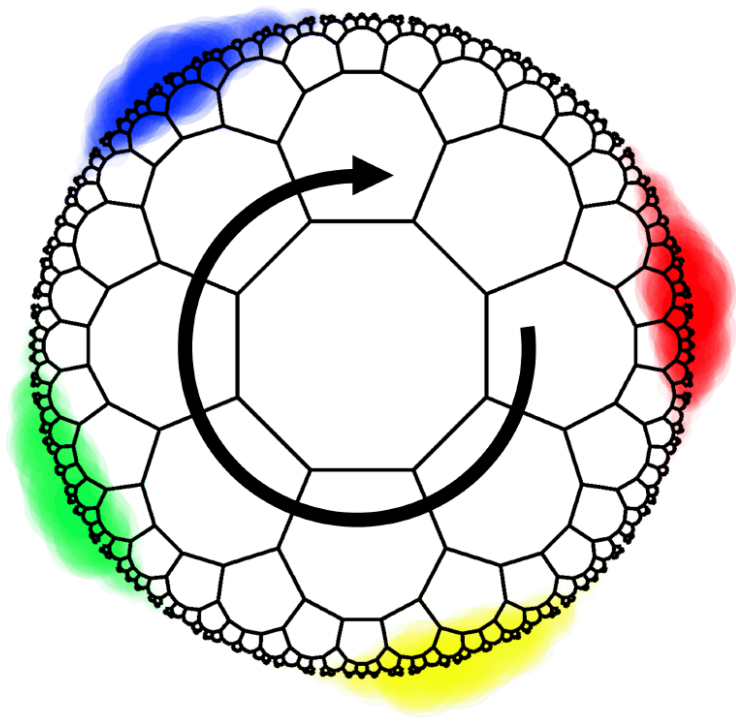
David M. Urwyler,¹ Patrick M. Lenggenger ^{1,2,3} Igor Boettcher ^{4,5}
Ronny Thomale ⁶ Titus Neupert ¹ and Tomáš Bzdušek ^{2,1,*}

arXiv:2203.07292 (th)

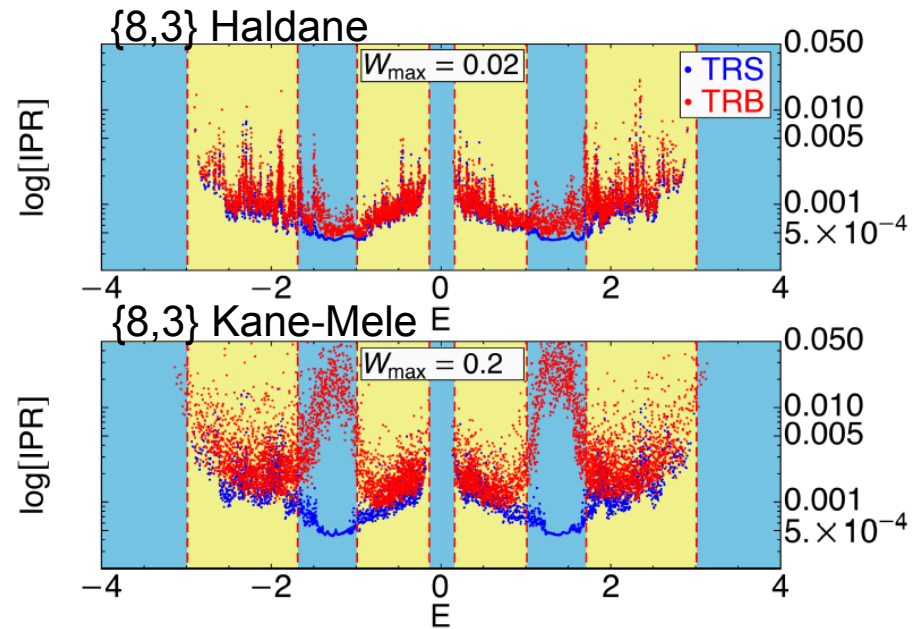
Hyperbolic Haldane model



Hyperbolic topological insulators



chiral/helical edge propagation

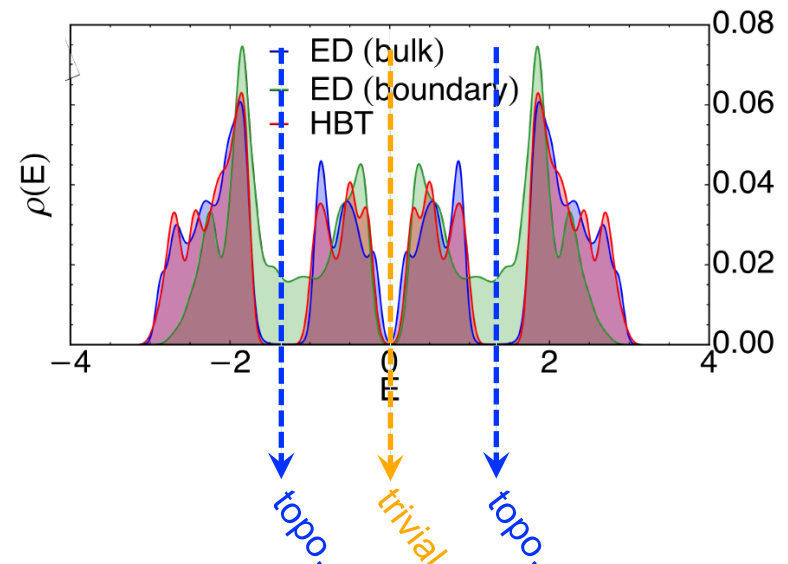
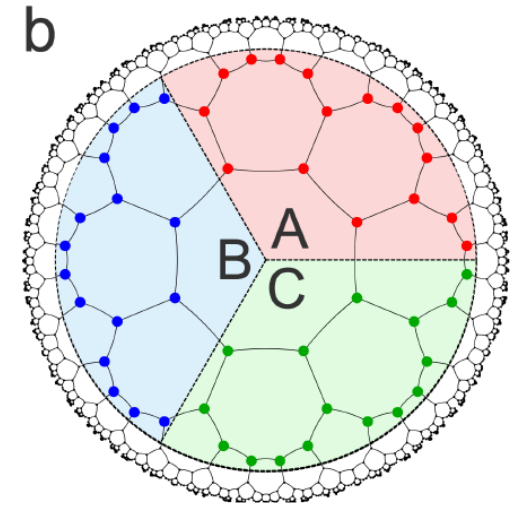


topological protection

Hyperbolic band topology

a

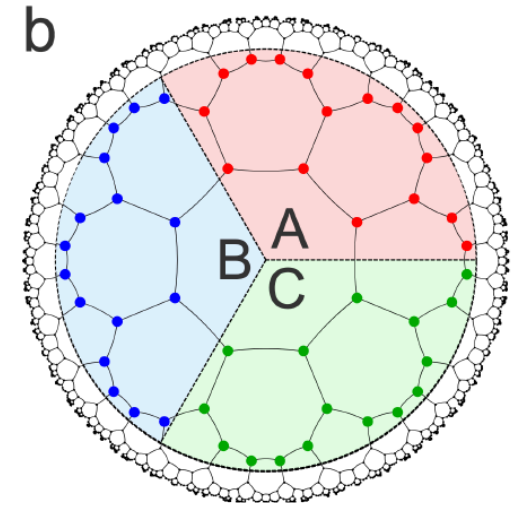
		Haldane		Kane-Mele	
f	μ		\mathcal{C}_{RS}		ν_{RS}
5/16	-1.3		-0.986		-0.971
8/16	0		0		0
11/16	+1.3		-0.986		-0.971



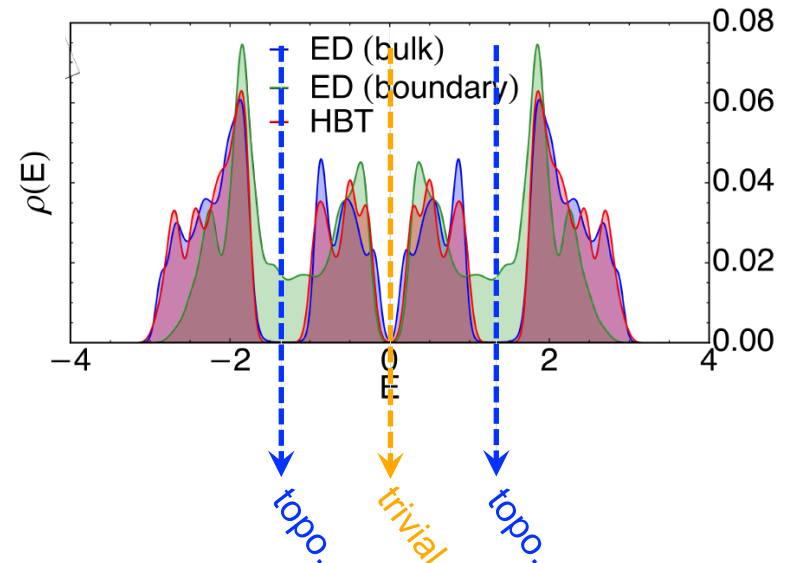
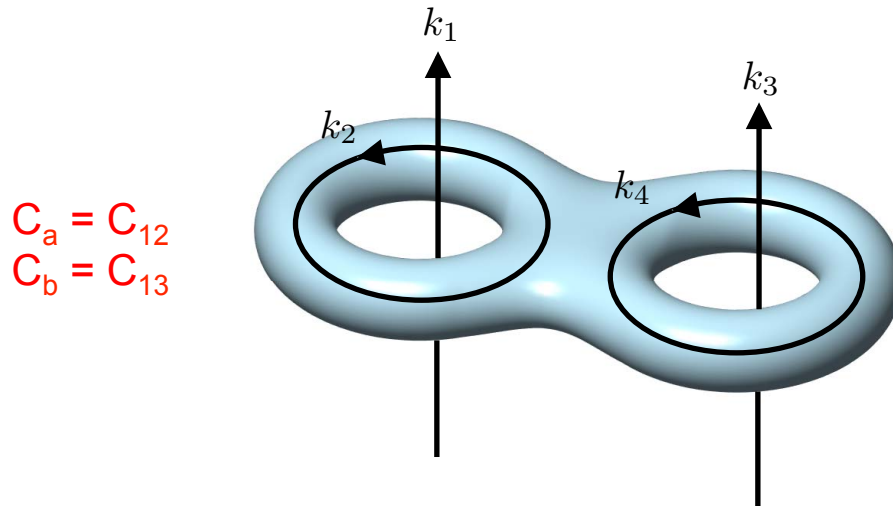
Hyperbolic band topology

a

f	μ	Haldane			Kane-Mele		
		C_a	C_b	C_{RS}	ν_a	ν_b	ν_{RS}
5/16	-1.3	-1	1	-0.986	1	1	-0.971
8/16	0	0	0	0	0	0	0
11/16	+1.3	-1	1	-0.986	1	1	-0.971



topological invariants in 2D planes of 4D Brillouin zone



Summary

- Hyperbolic lattices have a nonabelian translation group: $\Gamma \cong \pi_1(\Sigma_g)$
- Hyperbolic (automorphic) Bloch states: $\psi_{\mathbf{k}}(\gamma(z)) = \chi_{\mathbf{k}}(\gamma)\psi_{\mathbf{k}}(z)$
- Hyperbolic Brillouin zones: $\text{Jac}(\Sigma_g) \cong T^{2g} \cong \mathcal{M}(\Sigma_g, U(1))$
 $\mathcal{M}(\Sigma_g, U(r)), r > 1$
- Finite clusters with PBC correspond to $\Gamma_{\text{PBC}} \triangleleft \Gamma$
- Automorphic Bloch theorems (abelian/nonabelian)
- Hyperbolic flat bands
- Hyperbolic topological bands

Outlook

- Why does abelian HBT work so well?
- Parametrization of nonabelian Brillouin zones
- Physics of nonabelian Bloch states?
- Bulk-boundary correspondence ($C_{\text{real-space}} \leftrightarrow C_{k\text{-space}}$)?
- Implement PBC in experiment (e.g. electrical circuits)

JM & S. Rayan, *Sci. Adv.* 7, eabe9170 (2021)

JM & S. Rayan, *PNAS* 119, e2116869119 (2022)

I. Boettcher *et al.*, *PRB* 105, 125118 (2022)

Thank you!