Tensionless Strings Limits in 4d Conformal Manifolds

José Calderón Infante

Based on ongoing work with Irene Valenzuela

Swamplandia 2024, Seeon Abbey, 29/05/2024







Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

 $M_{tower} \sim e^{-\alpha \Delta \phi}$ as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$



Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

$$_{ver} \sim e^{-\alpha \,\Delta \phi}$$
 as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$

Precise order one bounds on the exponential rates • Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22] • Species scale Λ_{QG} : [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23]

• A pattern connecting them: [Castellano, Ruiz, Valenzuela '23]





Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

$$_{ver} \sim e^{-\alpha \,\Delta \phi}$$
 as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$

Precise order one bounds on the exponential rates • Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22] • Species scale Λ_{OG} : [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23] A pattern connecting them: [Castellano, Ruiz, Valenzuela '23]

Progress:





[Perlmutter, Rastelli, Vafa, Valenzuela '20]

Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

$$_{ver} \sim e^{-\alpha \,\Delta \phi}$$
 as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$

Precise order one bounds on the exponential rates • Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22] • Species scale Λ_{OG} : [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23] A pattern connecting them: [Castellano, Ruiz, Valenzuela '23]

Progress:





[Perlmutter, Rastelli, Vafa, Valenzuela '20]

Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

$$_{ver} \sim e^{-\alpha \, \Delta \phi}$$
 as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$

Precise order one bounds on the exponential rates • Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22] • Species scale Λ_{OG} : [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23] A pattern connecting them: [Castellano, Ruiz, Valenzuela '23]

Progress:

+ Bottom-up motivations

[Hamada, Montero, Vafa, Valenzuela '21] [Stout '21+'22] [JCI, Castellano, Herráez, Ibáñez '23]





[Perlmutter, Rastelli, Vafa, Valenzuela '20]

Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

$$_{ver} \sim e^{-\alpha \, \Delta \phi}$$
 as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$

Precise order one bounds on the exponential rates • Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22] • Species scale Λ_{OG} : [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23] A pattern connecting them: [Castellano, Ruiz, Valenzuela '23]

Progress:

+ Bottom-up motivations

[Hamada, Montero, Vafa, Valenzuela '21] [Stout '21+'22]

- [JCI, Castellano, Herráez, Ibáñez '23]
- + connections, pheno,





Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

$$_{ver} \sim e^{-\alpha \, \Delta \phi}$$
 as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$

Precise order one bounds on the exponential rates • Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22] • Species scale Λ_{OG} : [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23] A pattern connecting them: [Castellano, Ruiz, Valenzuela '23]

Progress:

+ Bottom-up motivations

[Hamada, Montero, Vafa, Valenzuela '21] [Stout '21+'22]

- [JCI, Castellano, Herráez, Ibáñez '23]
- + connections, pheno,



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

AdS/CFT basics: AdS CFT $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$

[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

AdS/CFT basics: CFT AdS $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

AdS/CFT basics: CFT AdS $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$

 $(\mathcal{M}, G_{ii}) \longleftrightarrow (\mathcal{M}_{CFT}, \chi_{ii})$



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

AdS/CFT basics: CFT AdS $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$

 $(\mathcal{M}_{CFT},\chi_{ii})$

Moduli space metric

$$\mathcal{L} \supset M_{Pl}^{d-1} \frac{1}{2} G_{ij} \partial_{\mu} \phi^{i} \partial^{\mu} \phi^{j}$$



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**





Moduli space metric

[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT ?

AdS/CFT basics: AdS CFT $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$ At infinite distance: Tower of operators with $\Delta - \Delta_{unitarity} \sim e^{-\alpha_{CFT}t}$



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

AdS/CFT basics: AdS CFT $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$

At infinite distance: Tower of operators with

 $\Delta - \Delta_{unitarity} \sim e^{-\alpha_{CFT}t}$

Question: Which operators? (e.g. unitarity bound depend on spin!)



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

AdS/CFT basics: AdS CFT $(\phi, m) \longleftrightarrow (\mathcal{O}, \Delta)$ At infinite distance: Tower of operators with $\Delta - \Delta_{unitarity} \sim e^{-\alpha_{CFT}t}$ **Question:** Which operators? (e.g. unitarity bound depend on spin!)

Higher-spin operators



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2





[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2



Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

■. Infinite distance → HS point

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

■. Infinite distance → HS point

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:**

Conformal manifold of local CFT in d>2

I. HS point → Infinite distance

II. Infinite distance → HS point

 $\Pi \cdot \gamma_{\ell} = \Delta_{\ell} - (\ell + d - 2) \sim e^{-\alpha_{\ell} t}$

Zamolodchikov distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:** Conformal manifold of local CFT in d>2 **I.** HS point \longrightarrow Infinite distance II. Infinite distance → HS point

$$\prod \gamma_{\ell} = \Delta_{\ell} - (\ell + d - 2) \sim e^{-\alpha_{\ell} t}$$

Zamolodchikov distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



+ existence of stress tensor is crucial!

[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:** Conformal manifold of local CFT in d>2 I. HS point → Infinite distance II. Infinite distance → HS point

$$\prod \gamma_{\ell} = \Delta_{\ell} - (\ell + d - 2) \sim e^{-\alpha_{\ell} t}$$

Zamolodchikov distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



Moreover:

Analogous statement for 2d CFTs



[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:** Conformal manifold of local CFT in d>2 I. HS point → Infinite distance II. Infinite distance → HS point

$$\prod \gamma_{\ell} = \Delta_{\ell} - (\ell + d - 2) \sim e^{-\alpha_{\ell} t}$$

Zamolodchikov distance

Local CFT: Posses stress tensor

Dynamical gravity in the bulk!



Today: Stringy origin of HS points **?** [JCI, Valenzuela '24]



Strings in the Conformal Manifold



- Inspiration: Emergent String Conjecture [Lee, Lerche, Weigand '19]
 - KK modes \rightarrow Decompactification
 - Excitations of weakly-coupled string

Strings in the Conformal Manifold



KK tower \rightarrow No HS fields



- Inspiration: Emergent String Conjecture [Lee, Lerche, Weigand '19]
 - KK modes \rightarrow Decompactification
 - Excitations of weakly-coupled string
 - String tower \rightarrow HS fields
 - **Expectation:** HS point \leftrightarrow tensionless string

Strings in the Conformal Manifold



KK tower \rightarrow No HS fields



Problem: $T_s \lesssim R_{AdS}^{-2} \longrightarrow$ String in a highly-curved background... hard to study!



- Inspiration: Emergent String Conjecture [Lee, Lerche, Weigand '19]

 - ✓ KK modes → Decompactification
 ✓ Excitations of weakly-coupled string
 - String tower \rightarrow HS fields
 - **Expectation:** HS point \leftrightarrow tensionless string

- Rely on CFT results and extract clues

A Distance Conjecture Approach

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

A Distance Conjecture Approach

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow \begin{array}{l} \text{Decompactific} \\ n \text{ extra dime} \end{array}$$


In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow \frac{\text{Decc}}{n \text{ ex}}$$



Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow$$

Decompactification of *n* extra dimensions

- $\alpha = \frac{1}{\sqrt{d-2}} \longrightarrow$ Emergent string limit
- **Caveat:** Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]
- From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

- Decompactification of $\alpha = \frac{1}{\sqrt{d-2}} \rightarrow \frac{1}{\text{string limit}}$
- Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

Three different values: $\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}}, \frac{1}{\sqrt{2}} \right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N



- Decompactification of $\alpha = \frac{1}{\sqrt{d-2}} \rightarrow \frac{1}{\text{string limit}}$
- Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N



- Decompactification of $\alpha = \frac{1}{\sqrt{d-2}} \rightarrow \frac{1}{\text{string limit}}$
- Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N



But...
$$\alpha \neq \frac{1}{\sqrt{3}}$$
 for all of them?

- Decompactification of $\alpha = \frac{1}{\sqrt{d-2}} \rightarrow \frac{1}{\text{string limit}}$
- Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow$$

Decompactificant *n* extra dimer

Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

Three different values:
$$\alpha = \begin{cases} \sqrt{\frac{2}{3}}, \sqrt{\frac{2}{3}}, \sqrt{\frac{2}{3}} \end{cases}$$

Out of 19 theories!

But... $\alpha \neq \frac{1}{\sqrt{3}}$ for all of them?

ation of
$$\alpha = \frac{1}{\sqrt{d-2}} \longrightarrow \frac{1}{\text{string limit}}$$

 $\left\{\frac{7}{12}, \frac{1}{\sqrt{2}}\right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] Suggests three different strings in AdS

Actually... Match $n = \{3,4,6\}$ \rightarrow Decompactification to $D = \{8,9,11\}$?

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow$$

Decompactificant *n* extra dimer

Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

Three different values:
$$\alpha = \begin{cases} \sqrt{\frac{2}{3}}, \sqrt{\frac{2}{3}}, \sqrt{\frac{2}{3}} \end{cases}$$

Out of 19 theories!

But... $\alpha \neq \frac{1}{\sqrt{3}}$ for all of them?

ation of
$$\alpha = \frac{1}{\sqrt{d-2}} \longrightarrow \frac{1}{\text{string limit}}$$

 $\left\{\frac{7}{12}, \frac{1}{\sqrt{2}}\right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] Suggests three different strings in AdS

Actually... Match $n = \{3,4,6\}$ \rightarrow Decompactification to $D = \{8,9,11\}$?

So... What is going on?!

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

- Decompactification of $\alpha = \frac{1}{\sqrt{d-2}} \rightarrow \frac{1}{\text{string limit}}$
- Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

Three different values: $\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}}, \frac{1}{\sqrt{2}} \right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N Three different values: $\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}}, \sqrt{\frac{1}{\sqrt{2}}} \right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] $\int \frac{1}{\sqrt{1-1}} E.g. \ \mathcal{N} = 4 \text{ SYM}$



In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

- **Caveat:** Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N Three different values: $\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] → E.g. $\mathcal{N} = 4$ SYM \checkmark Type IIB on AdS₅ × S⁵



Goal: Understand this case!

Type IIB on an 5-sphere

 $S = \frac{M_{Pl}^3}{2} \left[d^5 x \sqrt{-g} \left(R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R}) \right) \right]$

Type IIB on an 5-sphere $S = \frac{M_{Pl}^3}{2} \int d^5x \sqrt{-g} \left(R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R}) \right)$ Controls string coupling

Type IIB on an 5-sphere Controls 5-sphere radius $S = \frac{M_{Pl}^3}{2} \int d^5x \sqrt{-g} \left(R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R}) \right)$ Controls string coupling

Type IIB on an 5-sphere Controls 5-sphere radius $S = \frac{M_{Pl}^3}{2} \int d^5x \sqrt{-g} \left(R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R}) \right)$ Controls string coupling

KK tower

 $\frac{M_{KK}}{M_{Pl}} \sim e^{-\sqrt{\frac{8}{15}}\hat{R}}$

Type IIB on an 5-sphere Controls 5-sphere radius $S = \frac{M_{Pl}^3}{2} \int d^5x \sqrt{-g} \left(R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R}) \right)$ Controls string coupling

KK tower

 $\frac{M_{KK}}{M_{Pl}} \sim e^{-\sqrt{\frac{8}{15}}\hat{R}}$

$$R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R})$$

String tower







► Â











Convex Hull for N=4 SYM

 $\mathcal{N} = 4$ SU(N) gauge theory in 4d

Convex Hull for N=4 SYM

 $\mathcal{N} = 4$ SU(N) gauge theory in 4d



Convex Hull for N = 4 SYM

 $\mathcal{N} = 4$ SU(N) gauge theory in 4d





Convex Hull for N = 4 SYM

 $\mathcal{N} = 4$ SU(N) gauge theory in 4d









Problem:

No CFT distance in the N-direction :





Notice:

Convex hulls for AdS and CFT glue nicely together! (see later)



Notice:

Convex hulls for AdS and CFT glue nicely together! (see later)





KK tower \leftrightarrow **BPS operators**

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

 $M_{KK} \sim R_{AdS}^{-2\beta}$ • Weird S^5 stabilization

KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

Weird BPS $\frac{M_{KK} \sim R_{AdS}^{-2\beta} \longleftrightarrow \Delta_{BPS} \sim N^{\frac{2}{3}(1-2\beta)}}{\text{Weird } S^{5} \text{ stabilization}}$ spectrum



KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$







KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!

R

KK

AdS

Notice:

Convex hulls for AdS and CFT do not glue nicely together!



Anti-separation of scales: $\beta > 1/2 \rightarrow M_{KK} \ll R_{AdS}^{-1}$

't Hooft limit (fixed λ)



KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!

 \hat{R}

KK

AdS

Notice:

Convex hulls for AdS and CFT do not glue nicely together!

Weird BPS $M_{KK} \sim R_{AdS}^{-2\beta} \longleftrightarrow \Delta_{BPS} \sim N^{\frac{2}{3}(1-2\beta)}$ spectrum • Weird S^5 stabilization Long story short

Anti-separation of scales: $\beta > 1/2 \rightarrow M_{KK} \ll R_{AdS}^{-1}$

HS

't Hooft limit (fixed λ)

CFT




KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!



Weird BPS $M_{KK} \sim R_{AdS}^{-2\beta} \longleftrightarrow \Delta_{BPS} \sim N^{\frac{2}{3}(1-2\beta)}$ spectrum → Weird S^5 stabilization



KK tower \leftrightarrow **BPS operators**

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!



R

KK

Notice:

Convex hulls for AdS and CFT do not glue nicely together!

Weird BPS $M_{KK} \sim R_{AdS}^{-2\beta} \longleftrightarrow \Delta_{BPS} \sim N^{\frac{2}{3}(1-2\beta)}$ spectrum • Weird S^5 stabilization Long story short

Separation of scales: $\beta < 1/2 \rightarrow M_{KK} \gg R_{AdS}^{-1}$







KK tower \leftrightarrow **BPS operators**

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!



Notice:

Convex hulls for AdS and CFT do not glue nicely together!



KK tower \leftrightarrow **BPS operators**

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!



Notice:

Convex hulls for AdS and CFT do not glue nicely together!

KK tower ↔ BPS operators

Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} \sim R_{AdS}^{-1}$

No scale separation from the CFT!

R

KK

Notice:

Convex hulls for AdS and CFT do not glue nicely together!









Why $\alpha = \frac{1}{\sqrt{2}} \neq \frac{1}{\sqrt{3}}$ in $\mathcal{N} = 4$ SYM $\mathbf{?}$





$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 1:





Flat space emergent string:
$$\frac{M_s}{M_{Pl}} \rightarrow 0 + M_s \sim M_{KK}$$
 Here: $\frac{M_s}{M_{Pl}} \rightarrow 0 + M_s \ll M_{KK} \sim O(1)$

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 1:





$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$





$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$





$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Why $\alpha = \frac{1}{\sqrt{2}} \neq$

 $M_s \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

Why $\alpha = \frac{1}{\sqrt{2}} \neq$

What goes wrong when computing α ?

Recap

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

 $M_s \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!

1. Moduli space metric for g_s

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

 $M_s \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!

What goes wrong when computing α ?

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

- $M_s \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!
 - What goes wrong when computing α ?

Reason 2:

Recap

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

- $M_{S} \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!
 - What goes wrong when computing α ?
- 1. Moduli space metric for g_s 2. String excitation modes with g_s

Weakly curved: $M_s \sim \sqrt{T_s} \sim M_{Pl} g_s^{1/4}$

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

- $M_{s} \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!
 - What goes wrong when computing α ?

1. Moduli space metric for g_s 2. String excitation modes with g_s

Weakly curved: $M_s \sim \sqrt{T_s} \sim M_{Pl} g_s^{1/4}$ CFT: $M_s \sim M_{Pl} g_s^{1/2} \checkmark \neq$

CFT prediction for part of string spectrum in highly-curved AdS

Recap

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

- $M_{s} \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!
 - What goes wrong when computing α ?

1. Moduli space metric for g_s 2. String excitation modes with g_s

Weakly curved: $M_s \sim \sqrt{T_s} \sim M_{Pl} g_s^{1/4}$ CFT: $M_s \sim M_{Pl} g_s^{1/2} \checkmark \neq$

Weakly curved

CFT:

CFT predictio spectrum in highly-curved AdS

Recap

$$\neq \frac{1}{\sqrt{3}} \text{ in } \mathcal{N} = 4 \text{ SYM }$$

Reason 2:

- $M_{s} \ll R_{AdS}^{-1} \rightarrow$ Weakly curved approximation breaks down!
 - What goes wrong when computing α ?

1. Moduli space metric for g_s 2. String excitation modes with g_s

$$: M_{s} \sim \sqrt{T_{s}} \sim M_{Pl} g_{s}^{1/4}$$

$$M_{s} \sim M_{Pl} g_{s}^{1/2} \xrightarrow{4}$$

$$M_{s} \sim T_{s} R_{AdS}$$
Food for thought!
The string is a str

Recap: 4d SCFTs with simple gauge group (Lagrangian) admitting large N

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\} \quad \text{[Pe}$$

erlmutter, Rastelli, Vafa, Valenzuela '20]

Recap: 4d SCFTs with simple gauge group (Lagrangian) admitting large N

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\} \quad \text{[Pe}$$

New strings? Or same string, weirder background?

erlmutter, Rastelli, Vafa, Valenzuela '20]

Recap: 4d SCFTs with simple gauge group (Lagrangian) admitting large N

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\} \quad \text{[Pe}$$

New strings? Or same string, weirder background?

Problem: How to detect a string from the CFT?

erlmutter, Rastelli, Vafa, Valenzuela '20]

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\} \quad \text{[Pe}$$

New strings? Or same string, weirder background?

- **Recap:** 4d SCFTs with simple gauge group (Lagrangian) admitting large N
 - erlmutter, Rastelli, Vafa, Valenzuela '20]
 - g. $\mathcal{N} = 4$ SYM \checkmark Type IIB on AdS₅ × S⁵ \checkmark

 - **Problem:** How to detect a string from the CFT?
 - Instead, look for physical properties that are controlled only by α

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\} \quad \text{[Pe}$$

New strings? Or same string, weirder background?

- **Recap:** 4d SCFTs with simple gauge group (Lagrangian) admitting large N
 - erlmutter, Rastelli, Vafa, Valenzuela '20]
 - g. $\mathcal{N} = 4$ SYM \checkmark Type IIB on AdS₅ × S⁵

 - **Problem:** How to detect a string from the CFT?
 - Instead, look for physical properties that are controlled only by α
 - **1.** Ratio between *a* and *c* central charges

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}} \left(\frac{1}{\sqrt{2}} \right) \right\} \quad \text{[Pe}$$

New strings? Or same string, weirder background?

- **Recap:** 4d SCFTs with simple gauge group (Lagrangian) admitting large N
 - erlmutter, Rastelli, Vafa, Valenzuela '20]
 - g. $\mathcal{N} = 4$ SYM \checkmark Type IIB on AdS₅ × S⁵ \checkmark
 - **Problem:** How to detect a string from the CFT?
 - Instead, look for physical properties that are controlled only by α
 - **1.** Ratio between *a* and *c* central charges
 - 2. Hagedorn temperature at large N

$$\alpha = \sqrt{\frac{2c}{\dim G}}$$

$$\alpha = \sqrt{\frac{2c}{\dim G}} \quad \stackrel{\dim G = f(a,c)}{\longrightarrow} \alpha = \frac{1}{\sqrt{4\frac{a}{c} - 2}}$$

* Relevant for various aspects of low energy EFT!

* Relevant for various aspects of low energy EFT!

[Henningson, Skenderis '98]

Most notably: $a \neq c$ (at large N) \leftrightarrow No weakly-coupled Einstein gravity at low energies

$$m G = f(a, c)$$

$$\alpha = \frac{1}{\sqrt{4 \frac{a}{c} - 2}}$$

$$+ \frac{1}{4\alpha^2}$$
Depends
on α only

* Relevant for various aspects of low energy EFT!

```
[Henningson, Skenderis '98]
```

Most notably: $a \neq c$ (at large N) \leftrightarrow No weakly-coupled Einstein gravity at low energies

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE$$

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

 $\xrightarrow{T \to T_H} \infty \longleftarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies! \rightarrow **Expectation:** Hagedorn temperature should only depend on α

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

 $\xrightarrow{T \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$ **Hagedorn temperature:** $T_H \longrightarrow$ Controls exponential density of states at high energies! \rightarrow **Expectation:** Hagedorn temperature should only depend on α # chiral multiplets 4d $\mathcal{N} = 1$ SU(N) gauge theory \rightarrow 7 parameters: $\{n_{Ad}, n_F, n_{\bar{F}}, n_A, n_{\bar{A}}, n_{\bar{S}}, n_{\bar{S}}\}$
$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \xrightarrow{T \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$$
gedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!
$$\longrightarrow \text{Expectation: Hagedorn temperature should only depend on } \alpha \text{ !}$$

$$4d \mathcal{N} = 1 \text{ SU(N) gauge theory} \longrightarrow 7 \text{ parameters: } \left\{ n_{Ad}, n_F, n_{\bar{F}}, n_A, n_{\bar{A}}, n_S, n_{\bar{S}} \right\}$$
short... $Z(T) \to \infty \leftrightarrow \text{Hagedorn condition: } z_V(T_H) + \left\{ n_{Ad} + \frac{1}{2}(n_S + n_{\bar{S}} + n_A + n_{\bar{A}}) \right\} z_{\Phi}(T_H) = 1$

Hag

Long story

S

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \xrightarrow{T \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$$
gedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!
$$\implies \text{Expectation: Hagedorn temperature should only depend on } \alpha \text{ !}$$

$$= \text{I SU(N) gauge theory} \longrightarrow 7 \text{ parameters: } \left\{ n_{Ad}, n_F, n_{\bar{F}}, n_A, n_{\bar{A}}, n_S, n_{\bar{S}} \right\}$$
short... $Z(T) \to \infty \leftrightarrow \text{Hagedorn condition: } z_V(T_H) + \left\{ n_{Ad} + \frac{1}{2}(n_S + n_{\bar{S}} + n_A + n_{\bar{A}}) \right\} z_{\Phi}(T_H) = 1$

$$\mathcal{N} = 1 \text{ vector } \qquad \text{Nice... But not enough!}$$

Hag

Long story



$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \xrightarrow{T \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$$
gedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!
$$\implies \text{Expectation: Hagedorn temperature should only depend on } \alpha \text{ !}$$

$$= 1 \text{ SU(N) gauge theory} \longrightarrow 7 \text{ parameters: } \left\{ n_{Ad}, n_F, n_{\bar{F}}, n_A, n_{\bar{A}}, n_S, n_{\bar{S}} \right\}$$
short... $Z(T) \to \infty \leftrightarrow \text{Hagedorn condition: } z_V(T_H) + \left\{ n_{Ad} + \frac{1}{2}(n_S + n_{\bar{S}} + n_A + n_{\bar{A}}) \right\} z_{\Phi}(T_H) = 1$

$$\mathcal{N} = 1 \text{ vector} \qquad \text{Nice... But not enough!} \qquad \mathcal{N} = 1$$

Ha

Long story

(+) Anomaly cancellation + $\beta_{1-loop} = 0$: { \cdots } = $\beta - n_F$ Only 1 parameter!



$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \xrightarrow{T \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$$
gedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!
$$\implies \text{Expectation: Hagedorn temperature should only depend on } \alpha \text{ } \text{!}$$

$$4d \ \mathcal{N} = 1 \ \text{SU(N) gauge theory} \longrightarrow 7 \text{ parameters: } \left\{ n_{Ad}, n_F, n_{\bar{F}}, n_A, n_{\bar{A}}, n_S, n_{\bar{S}} \right\}$$
short... $Z(T) \to \infty \leftrightarrow \text{Hagedorn condition: } z_V(T_H) + \left\{ n_{Ad} + \frac{1}{2}(n_S + n_{\bar{S}} + n_A + n_{\bar{A}}) \right\} z_{\Phi}(T_H) = 1$

$$\mathcal{N} = 1 \text{ vector } \text{Nice... But not enough!} \qquad \mathcal{N} = 1$$

Ha

Long story

+ Anomaly cancellation + $p_{1-loop} = 0$: { · · · } = $3 - n_F$ Only 1 parameter!

Hagedorn condition: $z_V(T_H) + 3(3 - 4\alpha^2) z_{\Phi}(T_H) = 1$ Only depends on α !



Ha

Long story

 $n_F(\alpha)$

Expectation Hagedorn condition: $z_V(T_H) + 3(3 - 4\alpha^2) z_{\Phi}(T_H) = 1$ Only depends on α ! (magically) confirmed



$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

- $\xrightarrow{T \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$ **Hagedorn temperature:** $T_H \longrightarrow$ Controls exponential density of states at high energies!
 - \rightarrow **Expectation:** Hagedorn temperature should only depend on α
 - # chiral multiplets
 - 4d $\mathcal{N} = 1$ USp(2N)/SO(2N) gauge theory \rightarrow 3 parameters: $\{n_F, n_A, n_S\}$

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

Expectation: Hagedorn temperature should only depend on α

4d $\mathcal{N} = 1$ USp(2N)/SO(2N) gauge theory \rightarrow 3 parameters: $\{n_F, n_A, n_S\}$

Long story short... $Z(T) \to \infty \leftrightarrow$ Hagedorn condition: $z_V(T_H) + \{n_S + n_A\} z_{\Phi}(T_H) = 1$ $\mathcal{N} = 1$ vector \checkmark $\mathcal{N} = 1$ chirals

 $\xrightarrow{I' \to T_H} \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$

Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!

chiral multiplets

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

$$\bigoplus \beta_{1-loop} = 0: \{\cdots\}$$

 $\xrightarrow{I} \to T_H \longrightarrow \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$

- **Hagedorn temperature:** $T_H \longrightarrow$ Controls exponential density of states at high energies! \rightarrow **Expectation:** Hagedorn temperature should only depend on α # chiral multiplets 4d $\mathcal{N} = 1$ USp(2N)/SO(2N) gauge theory \rightarrow 3 parameters: $\{n_F, n_A, n_S\}$
 - Long story short... $Z(T) \to \infty \leftrightarrow$ Hagedorn condition: $z_V(T_H) + \{n_S + n_A\} z_{\Phi}(T_H) = 1$ $\mathcal{N} = 1$ vector \checkmark $\mathcal{N} = 1$ chirals
 - $= f(n_F)$ Only 1 parameter!

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies! \rightarrow **Expectation:** Hagedorn temperature should only depend on α # chiral multiplets 4d $\mathcal{N} = 1$ USp(2N)/SO(2N) gauge theory \rightarrow 3 parameters: $\{n_F, n_A, n_S\}$

$$\bigoplus \beta_{1-loop} = 0: \{\cdots\}$$



 $\xrightarrow{I} \to T_H \longrightarrow \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$

Long story short... $Z(T) \to \infty \leftrightarrow$ Hagedorn condition: $z_V(T_H) + \{n_S + n_A\} z_{\Phi}(T_H) = 1$ $\mathcal{N} = 1$ vector \checkmark $\mathcal{N} = 1$ chirals

 $= f(n_F)$ Only 1 parameter!

Hagedorn condition: $z_V(T_H) + 3(3 - 4\alpha^2) z_{\Phi}(T_H) = 1$ Only depends on α !

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

 \rightarrow **Expectation:** Hagedorn temperature should only depend on α

Long story short... $Z(T) \to \infty \leftrightarrow$ Hagedorn condition: $z_V(T_H) + \{n_S + n_A\} z_{\Phi}(T_H) = 1$ $\mathcal{N} = 1$ vector \checkmark $\mathcal{N} = 1$ chirals

$$(\mathbf{f}) \beta_{1-loop} = 0: \{ \cdots \}$$

 $n_F(\alpha)$

 $\xrightarrow{I} \to T_H \longrightarrow \infty \longrightarrow \rho(E) \sim e^{E/T_H} \text{ Stringy!}$

Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!

chiral multiplets 4d $\mathcal{N} = 1$ USp(2N)/SO(2N) gauge theory \rightarrow 3 parameters: $\{n_F, n_A, n_S\}$

 $= f(n_F)$ Only 1 parameter!

Expectation Hagedorn condition: $z_V(T_H) + 3(3 - 4\alpha^2) z_{\Phi}(T_H) = 1$ Only depends on α ! (magically) confirmed Same as for SU(N)



$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

Expectation: Hagedorn temperature should only depend on α **Confirmed**



Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

Expectation: Hagedorn temperature should only depend on α **Confirmed**

Caveat: Trouble with large numbers of flavors at large N



Hagedorn temperature: $T_H \longrightarrow$ Controls exponential density of states at high energies!

$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$

[Gadde, Pomoni, Rastelli '09] \rightarrow Restrict to flavor singlets!



$$Z(T) = \sum_{states} e^{-E/T} = \int \rho(E) e^{-E/T} dE \quad -$$



Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers



 $s \leftrightarrow \mathcal{N} = 2$ necklace quivers

 \rightarrow S¹ of orbifold singularities



Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

 S^1 of orbifold singularities

A very peculiar limit:

Driven by only axions \rightarrow Typically finite distance



Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

 S^1 of orbifold singularities

A very peculiar limit:

- Driven by only axions \rightarrow Typically finite distance
- **But!** CFT predicts infinite distance + HS conserved currents [Aharony, Berkooz, Rey '15]



Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

 S^1 of orbifold singularities

A very peculiar limit:

- Driven by only axions \rightarrow Typically finite distance
- **But!** CFT predicts infinite distance + HS conserved currents [Aharony, Berkooz, Rey '15]

Stringy origin?

Fundamental string remains tensionful...



Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

 S^1 of orbifold singularities

A very peculiar limit:

- Driven by only axions \rightarrow Typically finite distance
- **But!** CFT predicts infinite distance + HS conserved currents [Aharony, Berkooz, Rey '15]

Stringy origin?

- Fundamental string remains tensionful...
- D3 wrapping blow-up 2-cycle become tensionless! [Aharony, Berkooz, Rey '15]



String propagating in $AdS_5 \times S^1$!

Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

 S^1 of orbifold singularities

A very peculiar limit:

- Driven by only axions \rightarrow Typically finite distance
- **But!** CFT predicts infinite distance + HS conserved currents [Aharony, Berkooz, Rey '15]

Stringy origin?

- Fundamental string remains tensionful...
- D3 wrapping blow-up 2-cycle become tensionless! [Aharony, Berkooz, Rey '15]



String propagating in AdS₅ × S¹! Candidate for new emergent string in AdS \mathbf{Z} [Baume, JCI '20]

Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

 S^1 of orbifold singularities

A very peculiar limit:

- Driven by only axions \rightarrow Typically finite distance
- **But!** CFT predicts infinite distance + HS conserved currents [Aharony, Berkooz, Rey '15]

Stringy origin?

- Fundamental string remains tensionful...
- D3 wrapping blow-up 2-cycle become tensionless! [Aharony, Berkooz, Rey '15]

There is much to learn about/from the Distance Conjecture in AdS/CFT

There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side

There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side



There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side

Prove rest of CFT Distance Conjecture **?**

Distance in N-direction **?**

There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side

Prove rest of CFT Distance Conjecture **?**

Distance in N-direction **?**

Stringy side

There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side

Prove rest of CFT Distance Conjecture **?**

Distance in N-direction **?**

Stringy side



There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side

Prove rest of CFT Distance Conjecture **?**

Distance in N-direction **?**

Stringy side

New strings in AdS **?**

Building them: D3 wrapping blow-ups in AdS **?**

There is much to learn about/from the Distance Conjecture in AdS/CFT

CFT side

Prove rest of CFT Distance Conjecture **?**

Distance in N-direction **?**

Chank you for your attention!

Stringy side

New strings in AdS **?**

Building them: D3 wrapping blow-ups in AdS **?**