Axions in the dark dimension

Naomi Gendler Harvard University Swamplandia, May 29, 2024

based on 2404.15414 with Cumrun Vafa

Summary

Summary

A localized QCD axion in the dark dimension scenario is on the precipice of being detected experimentally.

Outline



1. The QCD axion and experiments



- 1. The QCD axion and experiments
- 2. The QCD axion in extra dimensions



- 1. The QCD axion and experiments
- 2. The QCD axion in extra dimensions
- 3. The QCD axion in the dark dimension

The QCD Lagrangian can, in principle, include a CP-violating term:

$$\mathcal{L} \supset \frac{\theta}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} \operatorname{tr} G_{\mu\nu} G_{\rho\sigma}$$

The QCD Lagrangian can, in principle, include a CP-violating term:

$$\mathcal{L} \supset \frac{\theta}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} \operatorname{tr} G_{\mu\nu} G_{\rho\sigma}$$

Experiments provide upper bounds on the neutron electric dipole moment, which show

$$\theta < 10^{-10}$$

The QCD Lagrangian can, in principle, include a CP-violating term:

$$\mathcal{L} \supset \frac{\theta}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} \operatorname{tr} G_{\mu\nu} G_{\rho\sigma}$$

Experiments provide upper bounds on the neutron electric dipole moment, which show

$$\theta < 10^{-10}$$

Strong CP problem: why is this number so small?

Promote θ to a dynamical field—an axion:

[Peccei, Quinn 1977]

$$\mathcal{L}_{\rm kin} = \frac{1}{2} f_a^2 \partial_\mu \theta \partial^\mu \theta$$

Promote θ to a dynamical field—an axion:

[Peccei, Quinn 1977]

$$\mathcal{L}_{\rm kin} = \frac{1}{2} f_a^2 \partial_\mu \theta \partial^\mu \theta$$

Non-perturbative QCD effects generate a potential for the axion:

$$V_{\rm QCD}(\theta) \sim \Lambda_{\rm QCD}^4 \theta^2 + \mathcal{O}(\theta^4)$$

Promote θ to a dynamical field—an axion:

[Peccei, Quinn 1977]

$$\mathcal{L}_{\rm kin} = \frac{1}{2} f_a^2 \partial_\mu \theta \partial^\mu \theta$$

Non-perturbative QCD effects generate a potential for the axion:

$$V_{\rm QCD}(\theta) \sim \Lambda_{\rm QCD}^4 \theta^2 + \mathcal{O}(\theta^4)$$

If this is the sole contribution to the QCD axion potential, then $\langle \theta \rangle = 0$ and the strong CP problem is solved: the axion dynamically relaxes the neutron EDM.

This field is called the "QCD axion" and it has mass



This field is called the "QCD axion" and it has mass



The QCD axion can couple to Standard Model fields.

This field is called the "QCD axion" and it has mass

$$m_a \approx \frac{\Lambda_{QCD}^2}{f_a}$$

The QCD axion can couple to Standard Model fields.

In particular, pions induce a coupling to photons:

$$\mathcal{L} \supset \frac{g^2 a}{16\pi^2 f_a} F \wedge F$$

This field is called the "QCD axion" and it has mass



The QCD axion can couple to Standard Model fields.

In particular, pions induce a coupling to photons:

$$\mathcal{L} \supset \frac{g^2 a}{16\pi^2 f_a} F \wedge F$$

Experiments are sensitive to the axion-photon coupling:

$$g_{a\gamma\gamma} = \frac{\alpha_{\rm EM}}{2\pi f_a}$$

Key idea: if axions exist and couple to photons, they will be produced in stars.

Key idea: if axions exist and couple to photons, they will be produced in stars.

If the axion-photon coupling is large enough, then stars cool too quickly.

Key idea: if axions exist and couple to photons, they will be produced in stars.

If the axion-photon coupling is large enough, then stars cool too quickly.

 \rightarrow this puts a bound on $g_{a\gamma\gamma}$

Key idea: if axions exist and couple to photons, they will be produced in stars.

If the axion-photon coupling is large enough, then stars cool too quickly.

 \rightarrow this puts a bound on $g_{a\gamma\gamma}$

Ensuring that neutron stars do not cool too quickly leads to the bound [Buschmann, Foster, Dessert, Long, Safdi 2021]

Key idea: if axions exist and couple to photons, they will be produced in stars.

If the axion-photon coupling is large enough, then stars cool too quickly.

 \rightarrow this puts a bound on $g_{a\gamma\gamma}$

Ensuring that neutron stars do not cool too quickly leads to the bound [Buschmann, Foster, Dessert, Long, Safdi 2021]

$$g_{a\gamma\gamma} \lesssim 3.3 \times 10^{-12} \text{ GeV}^{-1}$$

i.e. $f_a \gtrsim 3.5 \times 10^8 \text{ GeV}$

Direct detection experiments are also searching for axions.

Direct detection experiments are also searching for axions.

Axions produced in the sun can get converted to photons through $\mathcal{L}_{a\gamma\gamma} = g \, a \, \vec{E} \cdot \vec{B}$ coupling



Direct detection experiments are also searching for axions.

Axions produced in the sun can get converted to photons through $\mathcal{L}_{a\gamma\gamma} = g \, a \, \vec{E} \cdot \vec{B}$ coupling



The International Axion Observatory (IAXO) is the most sensitive upcoming experiment and will probe axions with decay constants

Direct detection experiments are also searching for axions.

Axions produced in the sun can get converted to photons through $\mathcal{L}_{a\gamma\gamma} = g \, a \, \vec{E} \cdot \vec{B}$ coupling



The International Axion Observatory (IAXO) is the most sensitive upcoming experiment and will probe axions with decay constants

 $f_a \sim 10^9 \text{ GeV}$

Let's consider a theory with:

Let's consider a theory with:

• 3 non-compact spatial dimensions

Let's consider a theory with:

- 3 non-compact spatial dimensions
- *n* mesoscopic spatial dimensions

Let's consider a theory with:

- 3 non-compact spatial dimensions
- *n* mesoscopic spatial dimensions
- *m* microscopic spatial dimensions
The QCD axion in extra dimensions

Let's consider a theory with:

- 3 non-compact spatial dimensions
- n mesoscopic spatial dimensions
- *m* microscopic spatial dimensions
- A localized Standard Model brane



The QCD axion in extra dimensions

Let's consider a theory with:

- 3 non-compact spatial dimensions
- *n* mesoscopic spatial dimensions
- *m* microscopic spatial dimensions
- A localized Standard Model brane



Let's further assume there exists a QCD axion.

The QCD axion in extra dimensions

Let's consider a theory with:

- 3 non-compact spatial dimensions
- *n* mesoscopic spatial dimensions
- *m* microscopic spatial dimensions
- A localized Standard Model brane



Let's further assume there exists a QCD axion.

This axion can propagate in 3 + p dimensions, with $p \ge 0$.

Swampland principles connect:

Swampland principles connect:

[Montero, Vafa, Valenzuela 2022]

size of the cosmological constant \leftrightarrow Kaluza-Klein scale

Swampland principles connect:

[Montero, Vafa, Valenzuela 2022]

size of the cosmological constant \leftrightarrow Kaluza-Klein scale



- 1 mesoscopic dimension
- $R \sim 1 10 \ \mu m$
- 5-dimensional Planck scale:

$$M_5 \sim 10^9 - 10^{10} \text{ GeV}$$

Swampland principles connect:

[Montero, Vafa, Valenzuela 2022]

size of the cosmological constant ↔ Kaluza-Klein scale



- 1 mesoscopic dimension
- $R \sim 1 10 \ \mu m$
- 5-dimensional Planck scale:

$$M_5 \sim 10^9 - 10^{10} \text{ GeV}$$

this is on the boundary of what is excluded by experiments!

Swampland principles connect:

[Montero, Vafa, Valenzuela 2022]

size of the cosmological constant \leftrightarrow Kaluza-Klein scale



- 1 mesoscopic dimension
- $R \sim 1 10 \ \mu m$
- 5-dimensional Planck scale:

$$M_5 \sim 10^9 - 10^{10} \text{ GeV}$$

this is on the boundary of what is excluded by experiments!

Our question: what do axion observations say about this scenario?

2 cases to consider.

2 cases to consider.

Case 1: QCD axion propagates in the fifth dimension



2 cases to consider.

Case 1: QCD axion propagates in the fifth dimension

Case 2: QCD axion is localized on the Standard Model brane



2 cases to consider.

Case 1: QCD axion propagates in the fifth dimension

Case 2: QCD axion is localized on the Standard Model brane

m m

Let's study each f_a in turn.



In the 5D theory, the axion has decay constant f_5 .



In the 5D theory, the axion has decay constant f_5 .

A reasonable expectation is that $f_5 \leq M_5$.



In the 5D theory, the axion has decay constant f_5 .

A reasonable expectation is that $f_5 \leq M_5$.

But the 4D decay constant is:

$$f_a^2 = f_5^3 R$$



In the 5D theory, the axion has decay constant f_5 .

A reasonable expectation is that $f_5 \lesssim M_5$.

But the 4D decay constant is:

$$f_a^2 = f_5^3 R$$

Which implies that

$$f_a \lesssim M_p$$



In the 5D theory, the axion has decay constant f_5 .

A reasonable expectation is that $f_5 \leq M_5$.

But the 4D decay constant is:

$$f_a^2 = f_5^3 R$$

Which implies that

$$f_a \lesssim M_p$$

so we don't get an interesting bound.

But there are KK copies of the axion!

But there are KK copies of the axion!

Each of them couples to photons, so one should ask if collective effects put rule out or constrain this scenario.

But there are KK copies of the axion!

Each of them couples to photons, so one should ask if collective effects put rule out or constrain this scenario.

$$g_{\mathrm{eff}} = \sqrt{\sum_{n \mid m_n \leq m_{\mathrm{exp}}} g_n^2}$$

But there are KK copies of the axion!

Each of them couples to photons, so one should ask if collective effects put rule out or constrain this scenario.

$$g_{ ext{eff}} = \sqrt{\sum_{n \mid m_n \leq m_{ ext{exp}}} g_n^2}$$

If $m_{\rm exp} > M_5$, then this sum is sensitive to $N \approx \frac{M_5}{M_{\rm KK}}$ modes

$$g_{\rm eff}^2 \approx \left(\frac{\alpha}{2\pi}\right)^2 \frac{M_5}{f_5^3}$$

But there are KK copies of the axion!

Each of them couples to photons, so one should ask if collective effects put rule out or constrain this scenario.

$$g_{ ext{eff}} = \sqrt{\sum_{n \mid m_n \leq m_{ ext{exp}}} g_n^2}$$

If $m_{\rm exp} > M_5$, then this sum is sensitive to $N \approx \frac{M_5}{M_{\rm KK}}$ modes $g_{\rm eff}^2 \approx \left(\frac{\alpha}{2\pi}\right)^2 \frac{M_5}{f_5^3}$

But for actual experiments $m_{\rm exp} \ll M_5$ and $g_{\rm eff}$ is not significantly enhanced.

But there are KK copies of the axion!

Each of them couples to photons, so one should ask if collective effects put rule out or constrain this scenario.

$$g_{ ext{eff}} = \sqrt{\sum_{n \mid m_n \leq m_{ ext{exp}}} g_n^2}$$

If $m_{\rm exp} > M_5$, then this sum is sensitive to $N \approx \frac{M_5}{M_{\rm KK}}$ modes $g_{\rm eff}^2 \approx \left(\frac{\alpha}{2\pi}\right)^2 \frac{M_5}{f_5^3}$

But for actual experiments $m_{\rm exp} \ll M_5$ and $g_{\rm eff}$ is not significantly enhanced.

Upshot: Bulk QCD axion in the dark dimension is not constrained.



Now let's consider the case where the QCD axion is localized on the Standard Model brane.



Now let's consider the case where the QCD axion is localized on the Standard Model brane.

This is what you would get if the Standard Model were realized on a blow-up cycle in a string theory compactification.



Now let's consider the case where the QCD axion is localized on the Standard Model brane.

This is what you would get if the Standard Model were realized on a blow-up cycle in a string theory compactification.



Argued in [Heckman, Vafa '08] that this is naturally follows from the hierarchy between $M_{\rm GUT}$ and $M_{\rm pl}$.

Now let's consider the case where the QCD axion is localized on the Standard Model brane.

This is what you would get if the Standard Model were realized on a blow-up cycle in a string theory compactification.



Argued in [Heckman, Vafa '08] that this is naturally follows from the hierarchy between $M_{\rm GUT}$ and $M_{\rm pl}$.

What is the 4-dimensional axion decay constant?



Again, a reasonable expectation is that $f_5 \leq M_5$.



Again, a reasonable expectation is that $f_5 \leq M_5$.

But what is f_a in the four-dimensional theory?



Again, a reasonable expectation is that $f_5 \leq M_5$.

But what is f_a in the four-dimensional theory?

We argue that for a *localized* axion, $f_a \leq M_5$.


Again, a reasonable expectation is that $f_5 \leq M_5$.

But what is f_a in the four-dimensional theory?

We argue that for a *localized* axion, $f_a \leq M_5$.

Easy way to see this: start with a bulk axion, with $S_5 \supset \frac{1}{2} \int d^5x \ f_5^3 (\nabla \theta)^2$



Again, a reasonable expectation is that $f_5 \leq M_5$.

But what is f_a in the four-dimensional theory?

We argue that for a *localized* axion, $f_a \leq M_5$.

Easy way to see this: start with a bulk axion, with

$$S_5 \supset \frac{1}{2} \int d^5 x \ f_5^3 (\nabla \theta)^2$$

Now suppose there exists a mechanism for localizing the axion.



Again, a reasonable expectation is that $f_5 \leq M_5$.

But what is f_a in the four-dimensional theory?

We argue that for a *localized* axion, $f_a \leq M_5$.

Easy way to see this: start with a bulk axion, with

$$S_5 \supset \frac{1}{2} \int d^5x \ f_5^3 (\nabla \theta)^2$$

Now suppose there exists a mechanism for localizing the axion.

Then we can integrate the axion field over the fifth dimension, with compact support on a region of size $1/M_5$.



Again, a reasonable expectation is that $f_5 \leq M_5$.

But what is f_a in the four-dimensional theory?

We argue that for a *localized* axion, $f_a \leq M_5$.

Easy way to see this: start with a bulk axion, with

$$S_5 \supset \frac{1}{2} \int d^5x \ f_5^3 (\nabla \theta)^2$$

Now suppose there exists a mechanism for localizing the axion.

Then we can integrate the axion field over the fifth dimension, with compact support on a region of size $1/M_5$.

$$f_a^2 \lesssim \frac{f_5^3}{M_5} \lesssim M_5^2$$
.

Consider a 5D theory with a large fifth dimension of size L.

Consider a 5D theory with a large fifth dimension of size $L\,.$

Consider a localize brane, supporting some U(1) gauge field.

Consider a 5D theory with a large fifth dimension of size $L\,.$

Consider a localize brane, supporting some U(1) gauge field.

Then particles localized on the brane should obey the WGC:

$$\frac{m}{M_p} \le g$$

Consider a 5D theory with a large fifth dimension of size $L\,.$

Consider a localize brane, supporting some U(1) gauge field.

Then particles localized on the brane should obey the WGC:

$$\frac{m}{M_p} \le g$$

Writing this inequality in terms of 5D quantities:

$$\frac{m}{M_5} \le g \ (M_5 L)^{\frac{1}{2}}$$

Consider a 5D theory with a large fifth dimension of size $L\,.$

Consider a localize brane, supporting some U(1) gauge field.

Then particles localized on the brane should obey the WGC:

$$\frac{m}{M_p} \le g$$

Writing this inequality in terms of 5D quantities:

$$\frac{m}{M_5} \le g \ (M_5 L)^{\frac{1}{2}}$$

m and *g* are independent of *L*, so we can consider this bound for arbitrary *L*. Taking $M_5 L = 1$, we find

$$\frac{m}{M_5} \le g.$$

Consider a 5D theory with a large fifth dimension of size $L\,.$

Consider a localize brane, supporting some U(1) gauge field.

Then particles localized on the brane should obey the WGC:

$$\frac{m}{M_p} \le g$$

Writing this inequality in terms of 5D quantities:

$$\frac{m}{M_5} \le g \ (M_5 L)^{\frac{1}{2}}$$

m and *g* are independent of *L*, so we can consider this bound for arbitrary *L*. Taking $M_5 L = 1$, we find

$$rac{m}{M_5} \le g.$$
 (for axions: $Sf \lesssim M_5$)

Consider a 5D theory with a large fifth dimension of size $L\,.$

Consider a localize brane, supporting some U(1) gauge field.

Then particles localized on the brane should obey the WGC:

$$\frac{m}{M_p} \le g$$

Writing this inequality in terms of 5D quantities:

$$\frac{m}{M_5} \le g \ (M_5 L)^{\frac{1}{2}}$$

m and *g* are independent of *L*, so we can consider this bound for arbitrary *L*. Taking $M_5 L = 1$, we find

$$\frac{m}{M_5} \le g.$$
 (for axions: $Sf \lesssim M_5$)

masses of localized particles are bounded by the higher dimensional Planck mass.

We've established that for the localized QCD axion,

 $f_a \lesssim M_5$.

We've established that for the localized QCD axion,

 $f_a \lesssim M_5$.

For the dark dimension scenario,

$$M_5 \sim 10^9 - 10^{10} \text{ GeV}$$
.

We've established that for the localized QCD axion,

 $f_a \lesssim M_5$.

For the dark dimension scenario,

$$M_5 \sim 10^9 - 10^{10} \text{ GeV}$$
.

But recall that experiments give us the bound

 $f_a \gtrsim 3.5 \times 10^8 \text{ GeV}$

We've established that for the localized QCD axion,

 $f_a \lesssim M_5$.

For the dark dimension scenario,

$$M_5 \sim 10^9 - 10^{10} \text{ GeV}$$
.

But recall that experiments give us the bound

 $f_a \gtrsim 3.5 \times 10^8 \text{ GeV}$

Gives a narrow window for the QCD axion: $m_a \sim (1-10) \text{ meV}$

 $10^8 \text{ GeV} \lesssim f_a \lesssim 10^9 - 10^{10} \text{ GeV}$

$$10^8 \text{ GeV} \lesssim f_a \lesssim 10^9 - 10^{10} \text{ GeV}$$

• This parameter regime will be probed by the next generation of solar axion experiments (IAXO).

$$10^8 \text{ GeV} \lesssim f_a \lesssim 10^9 - 10^{10} \text{ GeV}$$

• This parameter regime will be probed by the next generation of solar axion experiments (IAXO).

• New analyses of supernovae already start to cut into this region.

$$10^8 \text{ GeV} \lesssim f_a \lesssim 10^9 - 10^{10} \text{ GeV}$$

 This parameter regime will be probed by the next generation of solar axion experiments (IAXO).

• New analyses of supernovae already start to cut into this region.

• Misalignment dark matter production from this axion would be negligible – $\mathcal{O}(0.1 - 1\%)$ of the total dark matter abundance.

 Asked what bounds are placed on the dark dimension scenario if one assumes that there exists a QCD axion.

- Asked what bounds are placed on the dark dimension scenario if one assumes that there exists a QCD axion.
- Case 1: the bulk QCD axion lives in the same parameter space as the field theoretic axion.

- Asked what bounds are placed on the dark dimension scenario if one assumes that there exists a QCD axion.
- Case 1: the bulk QCD axion lives in the same parameter space as the field theoretic axion.
- Case 2: the localized QCD axion has a decay constant bounded by M_5 , which is a very strong constraint from the experimental perspective.

- Asked what bounds are placed on the dark dimension scenario if one assumes that there exists a QCD axion.
- Case 1: the bulk QCD axion lives in the same parameter space as the field theoretic axion.
- Case 2: the localized QCD axion has a decay constant bounded by M_5 , which is a very strong constraint from the experimental perspective.
- Upshot: a link between the size of a mesoscopic extra dimension, and the axion-photon coupling.