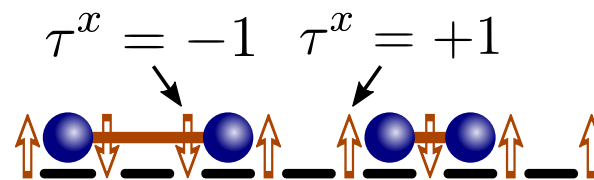


# Confinement in a one-dimensional $\mathbb{Z}_2$ lattice gauge theory at finite temperature



**Matjaž Kebrič, Jad C. Halimeh, Ulrich Schollwöck, and Fabian Grusdt**  
Munich Center for Quantum Science and Technology (MCQST) and  
Ludwig Maximilian University of Munich



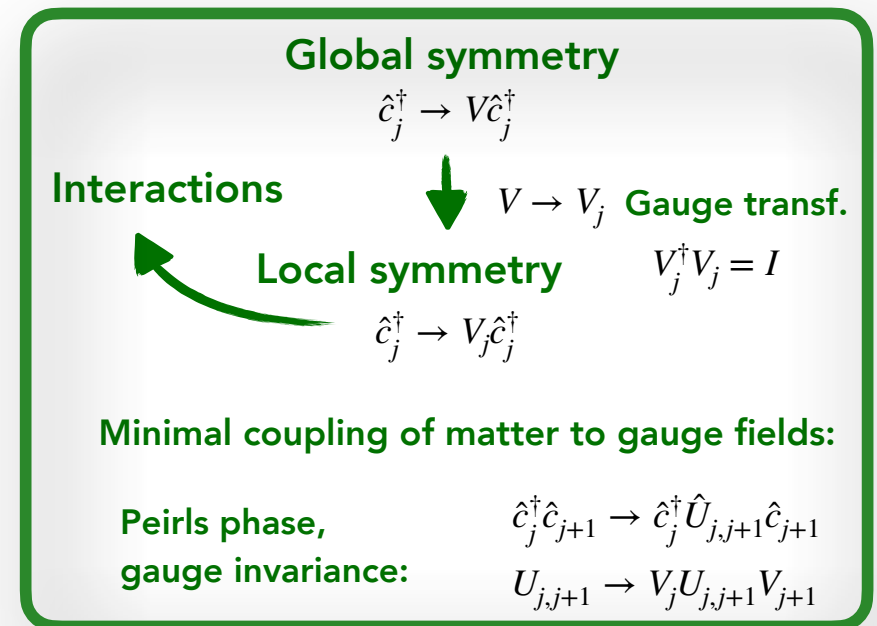
[www.quantummanybody.de](http://www.quantummanybody.de)

# Motivation

## Lattice gauge theories - "high energy physics"

- \* The standard model has local gauge symmetry:  $SU(3) \times SU(2) \times U(1)$ . Example:
  - ✓ Strong interaction: QCD -  $SU(3)$  gauge theory.
  - ✓ Electromagnetic interaction: QED -  $U(1)$  gauge theory.
- \* Non-perturbative problems: **lattice** gauge theory (LGT).
  - ✓ Analytics: lattice allows for regularization.
  - ✓ Numerics: Monte Carlo simulations **sign problem**
- \* Study of the **confinement-deconfinement** transition.
- \* Studying "simpler" toy models:  $\mathbb{Z}_N$  LGTs?

*Kogut, Rev. of Mod. Phys., 51, 659 (1979)*



*Zohar et al., Rep. Prog. Phys. 79 (2006).*



**Quantum simulation?!**

# Motivation

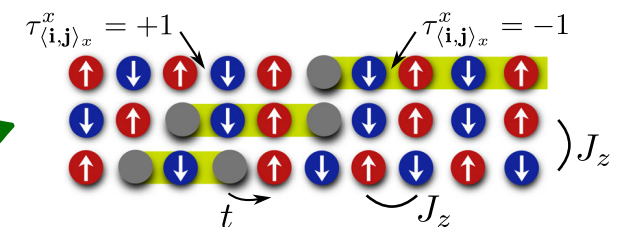
## Lattice gauge theories - condensed matter

- \* Study of the **confinement-deconfinement** transition.
- \* Emergent low energy theories in strongly correlated systems:
  - ✓ connection to high temperature superconductivity,
  - ✓ topological spin liquids,
  - ✓ fractional excitations, ...
- \* Direct mapping between a simple 1D  $\mathbb{Z}_2$  LGT and a  $t - J_z$  model.
- \* Many degrees of freedom: hard to solve!
  - ✓ Analytical and numerical calculations (DMRG, QMC).

Kogut, *Rev. of Mod. Phys.*, 51, 659 (1979)  
 Wilson, *PRD* 10, 2445 (1974)

Kitaev, *Ann. of Phys.* 303, 2 (2003).  
 Senthil et al., *PRB* 62, (2000).

Subir Sachdev *Rep. Prog. Phys.* 82 (2019).



Grusdt et al., *PRL* 125, (2020).

**Quantum simulation?!**

# Motivation

## Lattice gauge theories - quantum simulation

\* Many proposals:

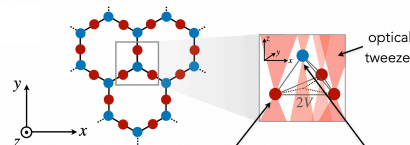
✓ Direct mapping between a simple 1D  $\mathbb{Z}_2$  LGT and a  $t - J_z$  model.

Spin-dependent super-lattice potential:

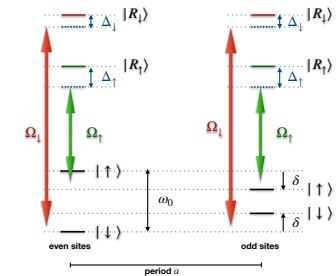
✓ Superconducting qubits,

✓ Rydberg tweezers, ...

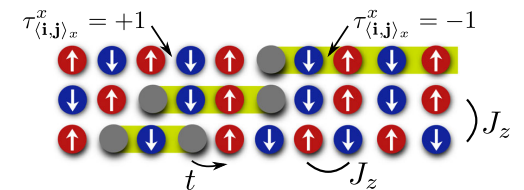
Homeier et al., PRB 104, (2021)



Homeier et al., Comm. Phys. 6, (2023).



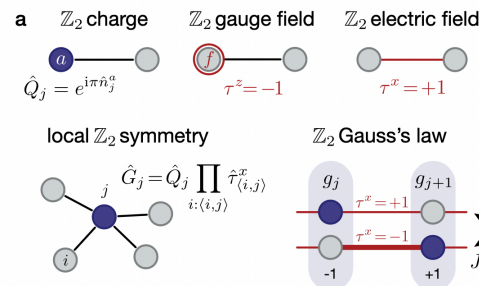
MK et al., PRL 127, (2021)



Grusdt et al., PRL 125, (2020).

\* Experimental realization using Floquet scheme.

Schweizer et al., Nature Phys. 15, (2019).



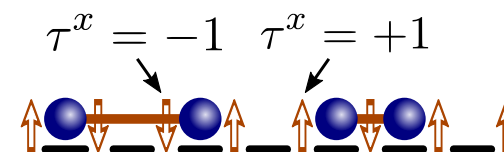
Görg et al., Nature Phys. 15, (2019).

Barbiero et al., Sci. Adv. 5 (2019)

# Overview of this talk

## 1D $\mathbb{Z}_2$ Lattice gauge theory coupled to dynamical matter

- \* Motivation
- \* The Hamiltonian and confinement of partons into mesons.
- \* Phase diagrams at different fillings.
- \* Confinement at finite temperature.
  - ✓ Smooth crossover from confined to thermally deconfined regime.
  - ✓ Simple experimental probe of confinement: Friedel oscillations and string length histograms.
  - ✓ Dynamical of confined mesons.
- \* Conclusion



*MK et al., Phys. Rev. B 109, (2024).*

# One-dimensional $\mathbb{Z}_2$ lattice gauge theory

## 1D $\mathbb{Z}_2$ lattice gauge theory

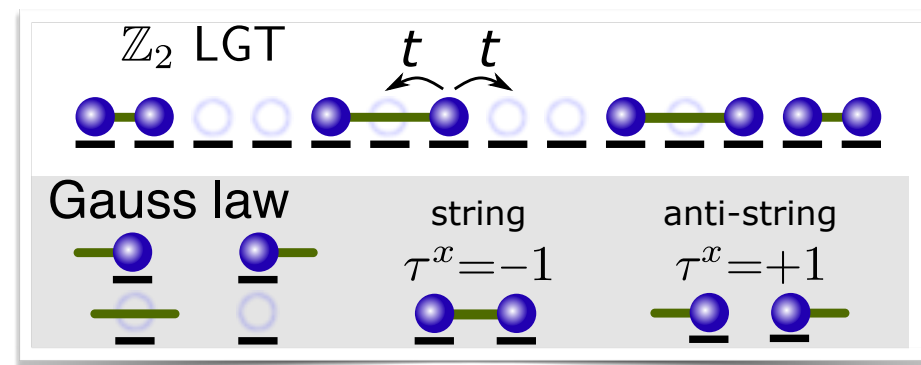
\*  $\hat{H} = -t \sum_{\langle i,j \rangle} (\hat{a}_i^\dagger \hat{\tau}_{\langle i,j \rangle}^z \hat{a}_j + \text{h.c.}) - h \sum_{\langle i,j \rangle} \hat{\tau}_{\langle i,j \rangle}^x + V \sum_j \hat{n}_j \hat{n}_{j+1}$

\* Gauss law:

$\hat{G}_i |\psi\rangle = \hat{\tau}_{\langle i-1,i \rangle}^x \hat{\tau}_{\langle i,i+1 \rangle}^x (-1)^{\hat{n}_i} |\psi\rangle = \pm |\psi\rangle \rightarrow g_i = +1$

$[\hat{H}, \hat{G}_j] = 0, \quad [\hat{G}_j, \hat{G}_i] = 0.$  *Prosko et al., PRB 96, (2017).*

*Borla et al., Phys. Rev. Lett. 124, (2020).*



*MK et al., Phys. Rev. Lett. 127, (2021).*

## Confinement

\* Linear confining potential  $h \sum_{\langle i,j \rangle} \hat{\tau}_{\langle i,j \rangle}^x$ .

\* Green's function decays exponentially for  $h > 0$   
*Borla et al., Phys. Rev. Lett. 124, (2020)*

$\rightarrow$  **Confinement of partons into mesons!**



## Solving the confinement problem

\* String-length basis: (b)  $\mathbb{Z}_2$  LGT string-length repr.

deconf.  $h = 0$   $\rightarrow$

conf.  $h \neq 0$   $\rightarrow$

*MK et al., Phys. Rev. Lett. 127, (2021).*

**Broken translational symmetry in the new basis  $\rightarrow$  confinement!**

# Ground state properties of the $\mathbb{Z}_2$ lattice gauge theory

## Phase Diagrams

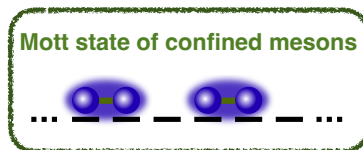
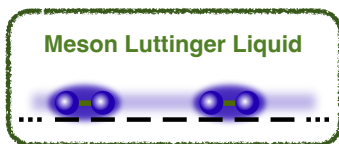
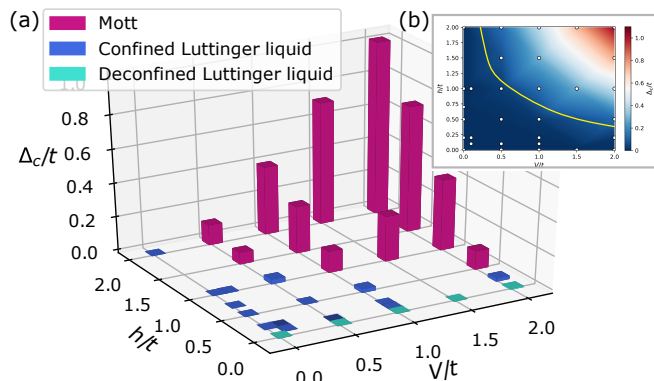
\* Luttinger liquid for generic values of  $h$  and  $V$ .

$$\hat{H} = -t \sum_{\langle i,j \rangle} \left( \hat{a}_i^\dagger \hat{\tau}_{\langle i,j \rangle}^z \hat{a}_j + \text{h.c.} \right) - h \sum_{\langle i,j \rangle} \hat{\tau}_{\langle i,j \rangle}^x + V \sum_j \hat{n}_j \hat{n}_{j+1}$$

\* Different Mott transitions at fillings  $n = 2/3$  and  $n = 1/2$ .

## Two-thirds filling

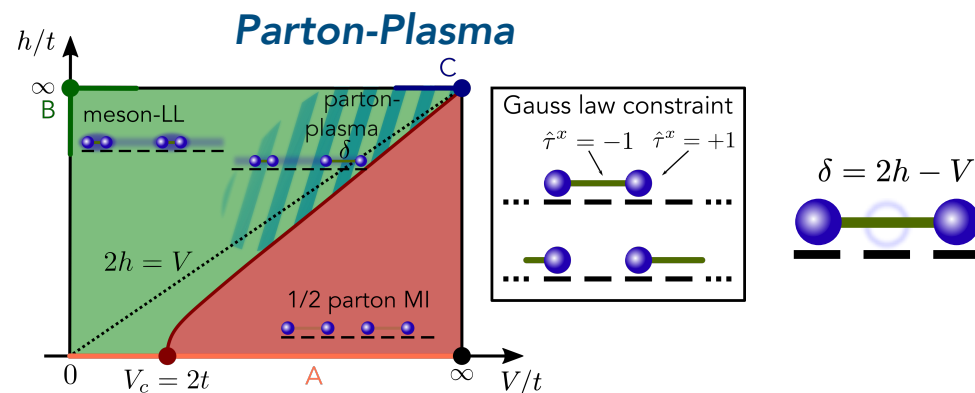
\* Mott state of confined dimers stabilized by  $V$ .



MK et al., Phys. Rev. Lett. 127, (2021)

## Half filling

\* Highly degenerate regime when  $h, V \gg t$  and  $2h \approx V$ .

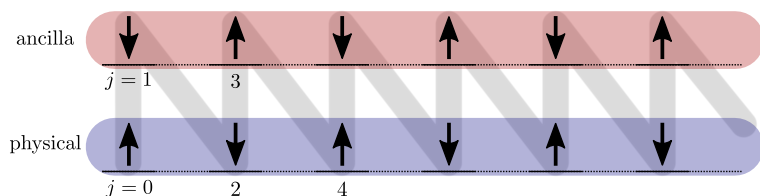


MK et al 2023 New J. Phys. 25 013035

# Numerical calculations of finite temperature LGT

## Method

- \* DMRG for the ground state calculations
- \* Quantum purification scheme to obtain finite temperature results.
- \* Ancilla site added to every physical lattice site.



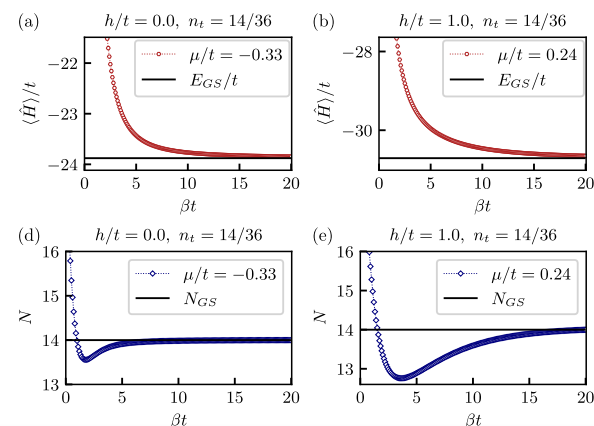
- \* Maximally entangled state between ancilla and physical lattice sites is the infinite temperature state  $\beta = 1/T = 0$

*Nocera et al., PRB, 93, 045 137, (2016).*

*Zwolak et al., PRL., 93, 207 205, (2004).*

*Feiguin et al., PRB, 72, 220 401, (2005).*

$$\hat{H} = -t \sum_{\langle i,j \rangle} \left( \hat{a}_i^\dagger \hat{\tau}_{\langle i,j \rangle}^z \hat{a}_j + \text{h.c.} \right) - h \sum_{\langle i,j \rangle} \hat{\tau}_{\langle i,j \rangle}^x$$



- \* By performing imaginary time evolution we "cool" the system

$$|\psi(\beta)\rangle = e^{-\beta \hat{\mathcal{H}}/2} |\psi(0)\rangle.$$

- \* Physical observables are calculated as

$$\langle \hat{O} \rangle = \frac{\langle \psi(\beta) | \hat{O} | \psi(\beta) \rangle}{\langle \psi(\beta) | \psi(\beta) \rangle}.$$



# Confinement in the $\mathbb{Z}_2$ LGT at finite temperature

## Green's function

- \*  $\mathbb{Z}_2$  invariant Green's function - probe of confinement:

$$\mathcal{G}(i-j) = \langle \hat{a}_i^\dagger \left( \prod_{i \leq l < j} \hat{\tau}_{l,l+1}^z \right) \hat{a}_j \rangle.$$

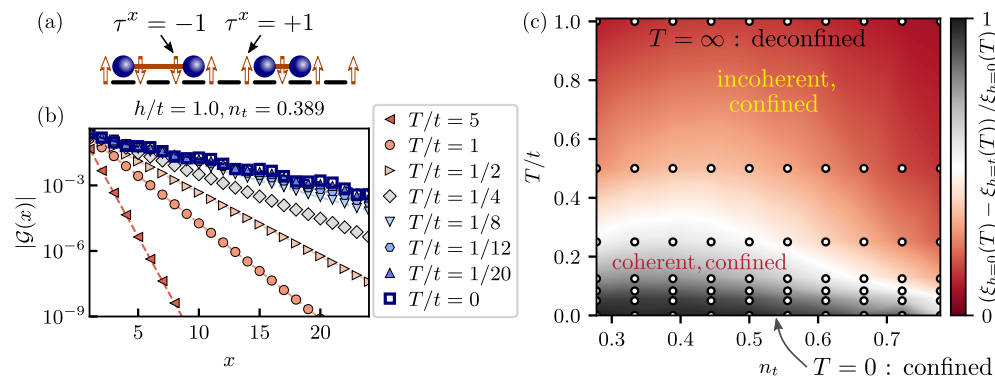
- \* Power-law decay in deconfined and exponential decay in the confined phase.

- \* We fit the Green's function with an exponential function to extract the correlation length  $\xi$

$$f_G(x) = Ax^{-\alpha} e^{-x/\xi}.$$

- \* Comparison of  $(\xi(h/t = 1) - \xi(h/t = 0)) / \xi(h/t = 0)$ .
- \* Quantum purification scheme via MPS.

$$\hat{H} = -t \sum_{\langle i,j \rangle} \left( \hat{a}_i^\dagger \hat{\tau}_{\langle i,j \rangle}^z \hat{a}_j + \text{h.c.} \right) - h \sum_{\langle i,j \rangle} \hat{\tau}_{\langle i,j \rangle}^x$$



MK et al., Phys. Rev. B 109, (2024).

## Smooth crossover! $T_c/t \approx 0.25$

- \* We uncover a smooth crossover region between confined and thermally deconfined regime as a function of temperature.

# 1+1D $\mathbb{Z}_2$ LGT at finite temperature

## Friedel oscillations

- \* Period of the Friedel oscillations doubles in the confined regime:  $2\pi n \rightarrow \pi n$ .

Borla et al., Phys. Rev. Lett. 124, (2020)

Parton Luttinger liquid  $h = 0$

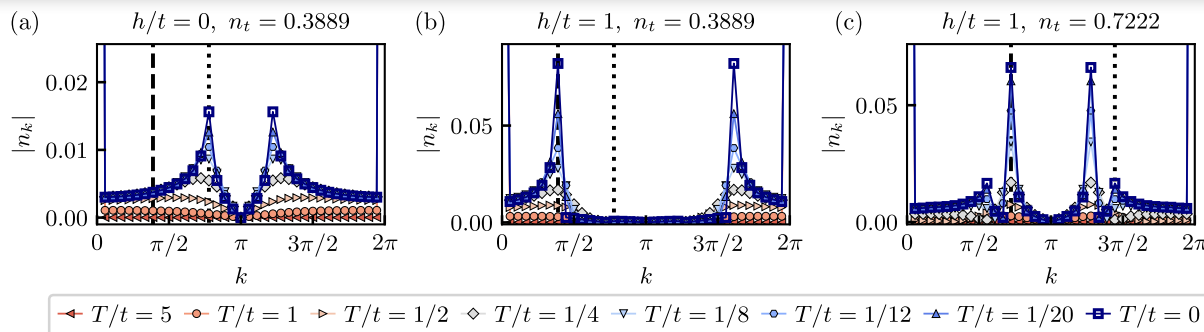
Meson Luttinger liquid  $h > 0$

Another signature of confinement!

$$\hat{H} = -t \sum_{\langle i,j \rangle} \left( \hat{a}_i^\dagger \hat{\tau}_{\langle i,j \rangle}^z \hat{a}_j + \text{h.c.} \right) - h \sum_{\langle i,j \rangle} \hat{\tau}_{\langle i,j \rangle}^x$$

- \* Frequency remains constant for  $h \neq 0$  regime as a function of  $T$ .

Pre-formed mesons at high temperature



Meson Luttinger,  $h > 0$

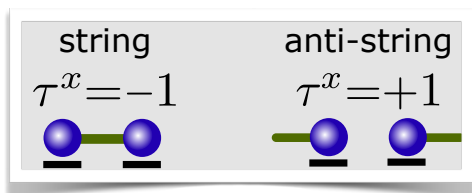


MK et al., Phys. Rev. B 109, (2024).

# 1+1D $\mathbb{Z}_2$ LGT at finite temperature

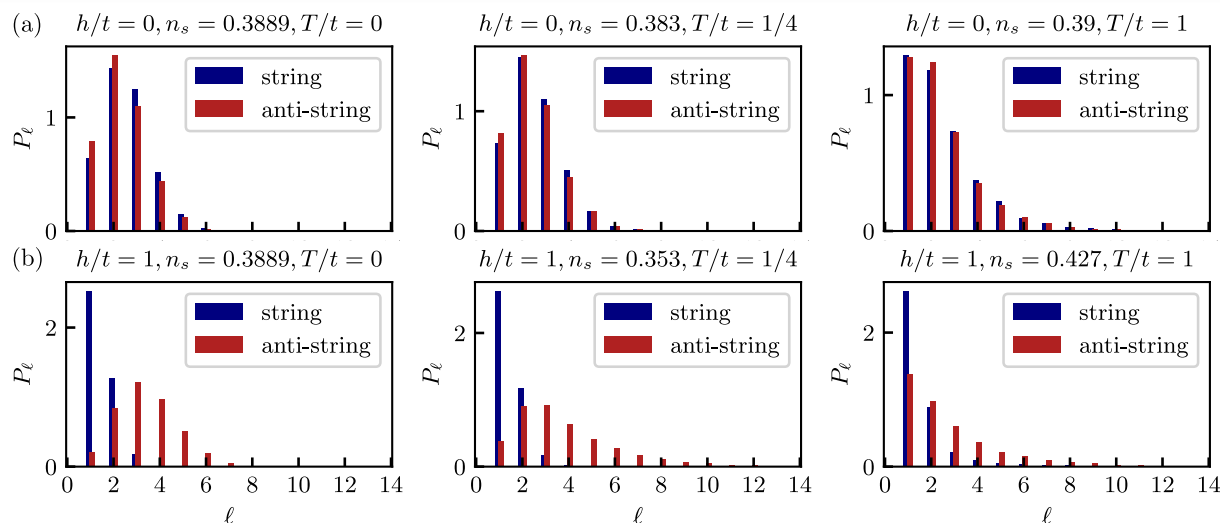
## Snapshots

- \* Green's function is difficult to access in cold-atom experiments.
- \* Snapshots sampled directly from MPS.
- \* String lengths can be easily accessible from projective measurements.
- \* Histograms of string and anti-string lengths.



**Bimodal distribution in the confined regime.**

**Robust measure of confinement!**

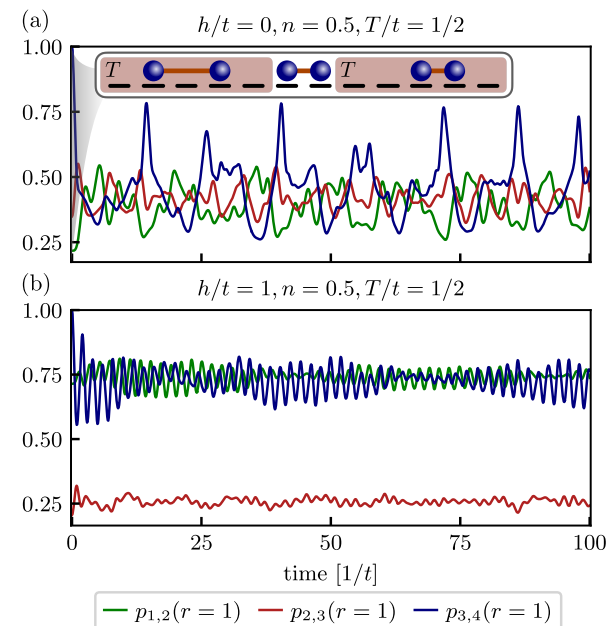
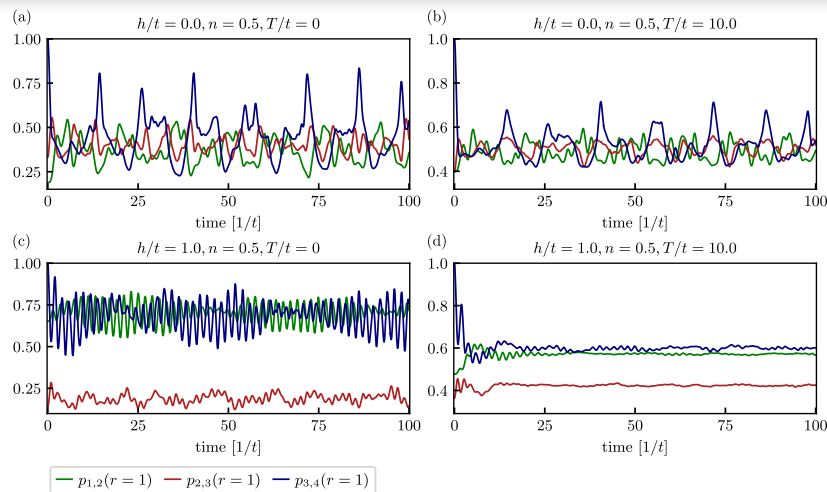


*MK et al., Phys. Rev. B 109, (2024).*

# Quench dynamics

## Time evolution

- \* Initial state with a meson in the center of the system.
- \* Probability of distance  $r = 1$  between charges.
- \* For  $h > 0$  the probability of well defined meson remains high at high temperature.



MK et al., Phys. Rev. B 109, (2024).

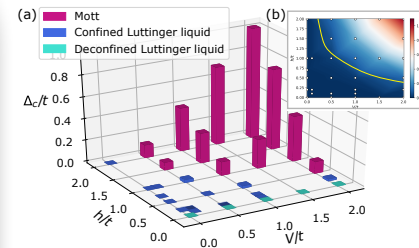
# Conclusion

## Summary

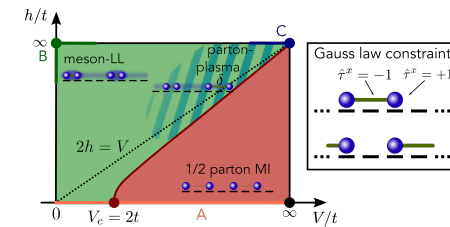
- \* Interplay of non-local confining field and local interaction result in rich phase diagrams.
- \* Crossover to deconfined regime at high temperature.
- \* Pre-formed mesons above the crossover temperature!
- \* Within the reach of current experimental setups.

## Outlook

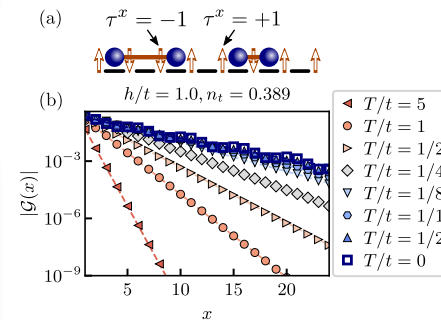
- \* Confinement in LGTs with more complicated gauge structure.
- \* Confinement in the 2D  $\mathbb{Z}_2$  LGT with matter?



MK et al., PRL, 127, (2021)



MK et al., New J. Phys. 25, (2023).



MK et al., Phys. Rev. B 109, (2024).

