Quantum computing for HEP: Digital twins for quantum states, some goals, and qudits

Christine Muschik

www.quantum-interactions.com

Let's see what your near future holds…

Overview

- 1. Digital twins for quantum states
- 2. Some QuantHEP goals
- 3. Qudits for the win

Neural shadow quantum state tomography

Towards digital twins for quantum states

Phys. Rev. Research 6, 023250 (2024).

Neural shadow quantum state tomography

Victor Wei, W. A. Coish, Pooya Ronagh, Christine A. Muschik Phys. Rev. Research 6, 023250 (2024).

Victor Wei **Stanford**

Bill Coish **McGill**

Pooya Ronagh IQC, PI, UWaterloo

Christine Muschik IQC, PI, UWaterloo

Digital twins of quantum states

Some quantum states are hard to find... **Example 20 and Some classically** Some classically.

Measure now **Measure now Measure now Measure now Measure in the me**

NISQ hardware **Post-process classically** Clean up the state: error mitigation

NISQ hardware **Post-process classically** Further optimization

Combine quantum computing with traditional LGT calculations

Finite density calculation, or time evolution **Example 20** ... regular lattice gauge theory computation.

Combine quantum computing with traditional LGT calculations

Finite density calculation, or time evolution **Example 20 million** ... regular lattice gauge theory computation.

Digital twins for quantum states

• Hybrid, etc…

-
- POVMs, etc…
-
- Other models…

Neural network quantum state tomography

Giacomo Torlai, Guglielmo Mazzola, Juan Carrasquilla, Matthias Troyer, Roger Melko & Giuseppe Carleo

Nature Physics 14, 447-450 (2018)

Neural network quantum state tomography

Neural quantum state

Transformer neural network quantum state ansatz

\n
$$
\sqrt{\frac{1}{\sqrt{1-\left(\frac{1}{2}\right)^{n}}\left(\frac{1}{2}\right)^{n}}}
$$
\nThese a bit-string $s = (s_1, s_2, \ldots, s_n) \in \{0, 1\}^n$

\nTake a bit-string $s = (s_1, s_2, \ldots, s_n) \in \{0, 1\}^n$

Parametrised by $\lambda = (\lambda_1, \lambda_2)$

Neural quantum state

 \blacksquare Measured bit-string *s* (corresponding to $|s\rangle$)

$$
\psi_{\lambda}(s) = \sqrt{p_{\lambda_1}(s)} e^{i\varphi_{\lambda_2}(s)}
$$

Complex-valued amplitude ⟨*s*|*ψλ*⟩

Pure state ansatz!

Г

Neural network quantum state tomography

Minimize: cross-entropy loss function

$$
L_{\lambda} = -\frac{1}{|\mathcal{B}|} \sum_{B \in \mathcal{B}} \sum_{s \in \{0,1\}^n} p_{\Phi}(s, B) \ln p_{\psi_{\lambda}}(s, B).
$$

Neural shadow quantum state tomography: NSQST

Classical shadows:

S. Aaronson, Shadow tomography of quantum states, in Proceedings of the 50th annual ACM SIGACT symposium on theory of computing (2018) pp. 325-338.

H.-Y. Huang, R. Kueng, and J. Preskill, Predicting many properties of a quantum system from very few measurements, Nature Physics 16, 1050 (2020).

Neural shadow quantum state tomography: NSQST

1. Use classical shadows 2. New cost function: fidelity

- 1. Improved phase information
- 2. Noise robustness

Neural shadow quantum state tomography: NSQST

- Perform random Clifford tails U_i and measure bit strings $|b_i\rangle$.
- Collect stabilizer states: $|\phi_i\rangle = U_i^{\dagger} |b_i\rangle$.
- Average effect of the Clifford twirling is a depolarizing noise channel M with strength $(2^{n} + 1)^{-1}$.
- Classical shadows: $\rho_i = \mathcal{M}^{-1}(|\phi_i\rangle\langle\phi_i|)$.
- Target state: $|\Phi\rangle\langle\Phi| = \mathbb{E}[\mathcal{M}^{-1}(|\phi_i\rangle\langle\phi_i|)].$
- New loss:

$$
1 - |\langle \psi_{\lambda} | \Phi \rangle| \approx 1 - \frac{1}{N} \sum_{i}^{N} Tr(O_{\lambda} \rho_{i}) = 1 - \frac{1}{2^{n}} \left(1 - \frac{1}{f} \right) - \frac{1}{N} \sum_{i}^{N} |\langle \phi_{i} | \psi_{\lambda} \rangle|^{2}
$$

NSQST for concrete examples

1D-QCD: time-evolved state

Heisenberg antiferromagnet

Phase shifted GHZ state

NSQST for a concrete example

Simulating one-dimensional quantum chromodynamics on a quantum computer: Real-time evolutions of tetra- and pentaquarks

Yasar Y. Atas*,^{1,2,}¹] Jan F. Haase*,^{1,2,3},¹] Jinglei Zhang,^{1,2},⁸] Victor Wei,^{1,4} Sieglinde M.-L. Pfaendler,⁵ Randy Lewis,⁶ and Christine A. Muschik^{1,2,7}

NSQST for a 1D-QCD time evolution

Infidelity

NSQST with pre-training

NSQST for a 1D-QCD time evolution

NSQST: Robustness to noise

Phys. Rev. Research 6, 023250 (2024).

Outlook

Food for thought

Quantum experiments are difficult to perform:

- Expensive
- Time-intensive

No-cloning theorem:

- How to capture experiments in a re-usable fashion
- Make the results of an experiment available to the community

Neural network trained on quantum data:

- Schemes like NSQST can be more useful than a list of measurements
- in the same way GPT is more useful than a text on the web

The neural network representation is a **digital twin** of the quantum state

- Shift the value from quantum experiments to **quantum data**
- Inference from the digital twin us **cheaper** than re-running quantum experiments
- The digital town is more **malleable and easier to interface with** than a quantum computer

Combine quantum computing with traditional LGT calculations

Finite density calculation, or time evolution **Example 20** ... regular lattice gauge theory computation.

Hamiltonian learning

Time evolution **Local measurements**
on time-evolved state

Matrix Product States

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Some goals

 $Fast \rightarrow Present \rightarrow Future$

Real-time dynamics of lattice gauge theories with a fewqubit quantum computer

Nature 534, 516-519 (2016)

Quantum simulation of 1D-QED

Theory + Experiment

Top 10 breakthroughs in physics 2016

Physics World

nature

COMPUTING

In a First, Quantum Computer **Simulates High-Energy Physics**

The technique could allow quantum computers to address otherwise-intractable problems in particle physics

Some goals

Gauge group of QCD

Loop-String-Hadron Formulation for SU(2) and SU(3): Indrakshi Raychowdhury, Jesse Stryker General quantum algorithms for non-Abelian LGTs: [Zohreh Davoudi,](https://arxiv.org/search/?searchtype=author&query=Davoudi%2C+Z) [Alexander F. Shaw,](https://arxiv.org/search/?searchtype=author&query=Shaw%2C+A+F) [Jesse R. Stryker](https://arxiv.org/search/?searchtype=author&query=Stryker%2C+J+R) Trailhead for SU(3): [Anthony Ciavarella,](https://arxiv.org/search/?searchtype=author&query=Ciavarella%2C+A) [Natalie Klco](https://arxiv.org/search/?searchtype=author&query=Klco%2C+N), [Martin J. Savage](https://arxiv.org/search/?searchtype=author&query=Savage%2C+M+J) Digitizing SU(2) Gauge Fields: T. Hartung, T. Jakobs, K. Jansen, J. Ostmeyer, C. Urbach 3 new quantum simulation ideas for non-Abelian gauge theories: Torsten Zache, Daniel González-Cuadra, Peter Zoller

Simulating 1D-QCD on a quantum computer

Theory + Experiment:

- Tetraquarks
- Real-time evolution

Simulate quarks with three colours, i.e. the gauge group of QCD

Simulating 1D-QCD on a quantum computer

Related work:

Preparations for Quantum Simulations of Quantum Chromodynamics in 1+1 Dimensions: (I) Axial Gauge

Roland C. Farrell, Ivan A. Chernyshev, Sarah J. M. Powell, Nikita A. Zemlevskiy, Marc Illa, Martin J. Savage

Tools necessary for quantum simulations of $1+1$ dimensional quantum chromodynamics are developed. When formulated in axial gauge and with two flavors of quarks, this system requires 12 qubits per spatial site with the gauge fields included via non-local interactions. Classical computations and D-Wave's quantum annealer Advantage are used to determine the hadronic spectrum, enabling a decomposition of the masses and a study of quark entanglement. Color edge states confined within a screening length of the end of the lattice are found. IBM's 7-qubit quantum computers, ibmq_jakarta and ibm_perth, are used to compute dynamics from the trivial vacuum in one-flavor QCD with one spatial site. More generally, the Hamiltonian and quantum circuits for time evolution of $1 + 1$ dimensional $SU(N_c)$ gauge theory with N_f flavors of quarks are developed, and the resource requirements for large-scale quantum simulations are estimated.

M. Savage, University of Washington

Experimental demonstration of color neutral objects (gauge singlets/color singlets).

Color neutral states of SU(3): invariant under arbitrary rotations in color space

Involve all the color charges available in the theory i.e. red, green, and blue (and their anticolors). In contrast to Abelian quantum electrodynamics (QED), where a singlet state involves electron-positron pairs only.

Color singlet states are the relevant physical states,

 \rightarrow important step towards the understanding, description and prediction of more complex and realistic experiments

Some goals

Some goals

Gauge theories for particle physics beyond 1D

Experimental demonstrations

Gauge theories for particle physics beyond 1D

- N. Klco, M. J. Savage, and J. R. Stryker, Phys. Rev. D. 101, 074512 (2020).
- A. Ciavarella, N. Klco, and M. J. Savage, Phys. Rev. D 103, 094501 (2021) .
- S. A Rahman, R. Lewis, E. Mendicelli, and S. Powell, Phys. Rev. D 104, 034501 (2021).
- A. N. Ciavarella and I. A. Chernyshev, Phys. Rev. D 105, 074504 (2022).
- S. A Rahman, R. Lewis, E. Mendicelli, and S. Powell, Phys. Rev. D 106, 074502 (2022).
- S. A. Rahman, R. Lewis, E. Mendicelli, and S. Powell, "Real time evolution and a traveling excitation in $SU(2)$ " pure gauge theory on a quantum computer," (2022) , $arXiv:2210.11606$ [hep-lat].
- A. N. Ciavarella, "Quantum Simulation of Lattice QCD with Improved Hamiltonians," (2023) , arXiv:2307.05593 [hep-lat].

Impressive Advances!

(so far gauge fields or matter fields were trivial)

Experimental demonstration

Gauge theories for particle physics beyond 1D

Including both - dynamical gauge and matter fields

Some goals

Representing the gauge fields

Gauge fields

Truncations for bosonic systems:

- Schwinger boson representation
- Holstein-Primakoff-representation
- Dysen-Maleev transformation
- Highly occupied boson model

Truncations for qubit systems:

$$
\hat{E} \longmapsto \hat{S}^z, \qquad \hat{V}^- \equiv \begin{bmatrix} 0 & \dots & \dots & 0 \\ 1 & \dots & \dots & 0 \\ 0 & \ddots & \vdots & 0 \\ 0 & \dots & 1 & 0 \end{bmatrix} \quad \begin{array}{l}\n \text{Alternative:} \\
\hat{V}^- \equiv \hat{S}^- / |l| \\
\hat{S}^- \equiv \hat{S}^x - i\hat{S}^y\n \end{array}
$$

Gauge fields

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\hat{S}^- \equiv \hat{S}^x - i\hat{S}^y\n \end{array}
$$

Hilbert space truncation with a regularisation of the gauge group

(*) Our scheme was further improved in: C. W. Bauer and D. M. Grabowska, Phys. Rev. D 107, L031503 (2023).

Hilbert space truncation with a regularisation of the gauge group

g = bare coupling

Number of states required to reach a 1% accuracy in the expectation value of the two-dimensional plaquette in QED

Some goals

Some goals

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Beyond binary

Today's quantum hardware: **capable of qudit encoding**

- Trapped ions
- Superconducting architectures
- Rydberg atoms in optical tweezers
- Ultracold atoms in optical lattices
- Nuclear spins
- Photonic systems

Gauge fields represented by qudits

[Reference to Martin's paper]

Qudits for Quantum Technology

Positions available (Masters/PhD/Postdoc)

www.quantum-interactions.com

Thank you for your time

Introduced "B-rep" Quantum 5, 393 (2021).

Developed qudit algorithms

1D-SU(2) Experiment on IBM 1D-SU(3) Experiment on IBM